

Partially Acoustic Dark Matter & Large Scale Structure

Yuhsin Tsai

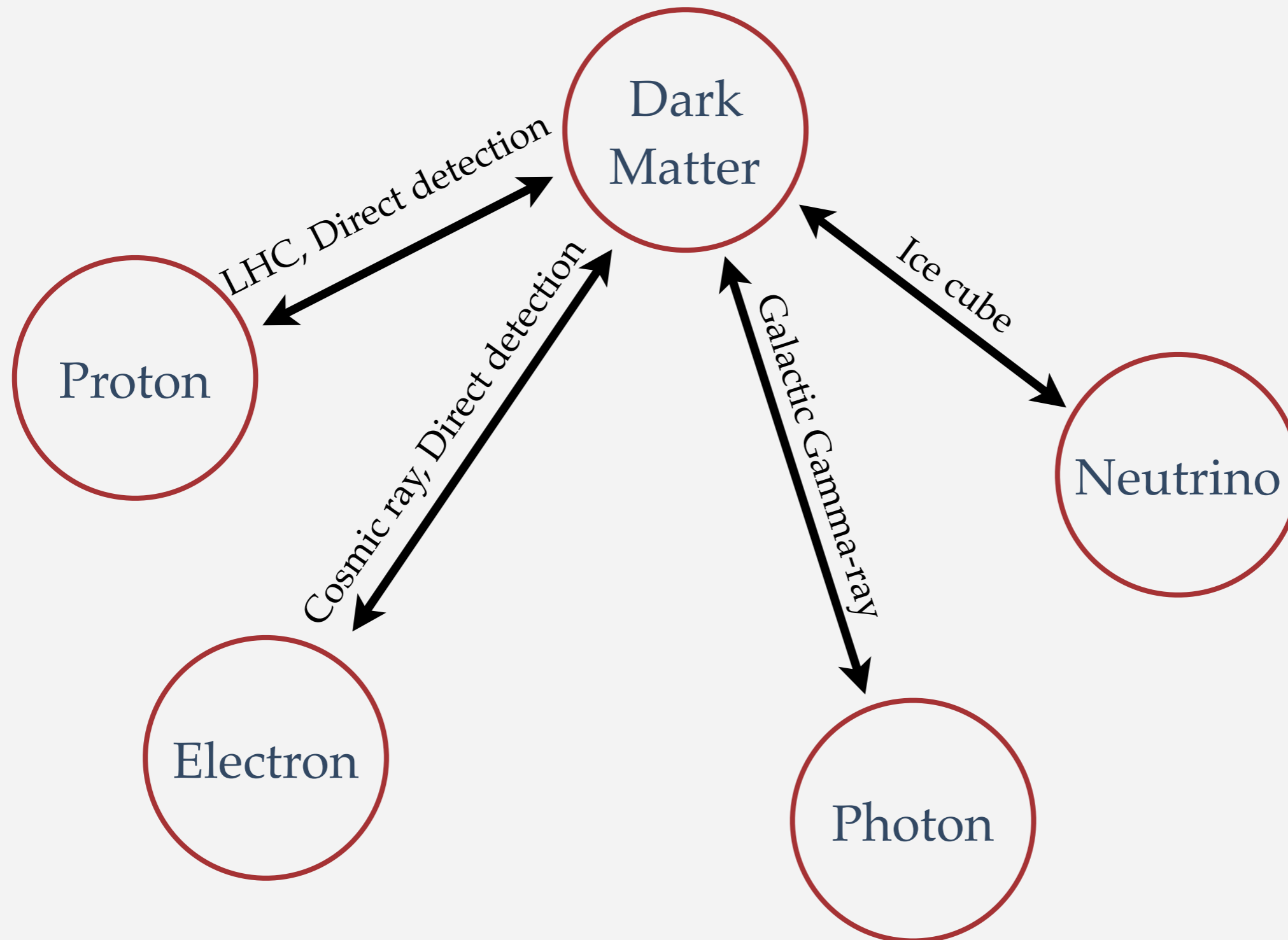
University of Maryland

JHEP1612(2016)108, Zackaria Chacko, Yanou Cui, Sungwoo Hong, and Takemichi Okui, YT

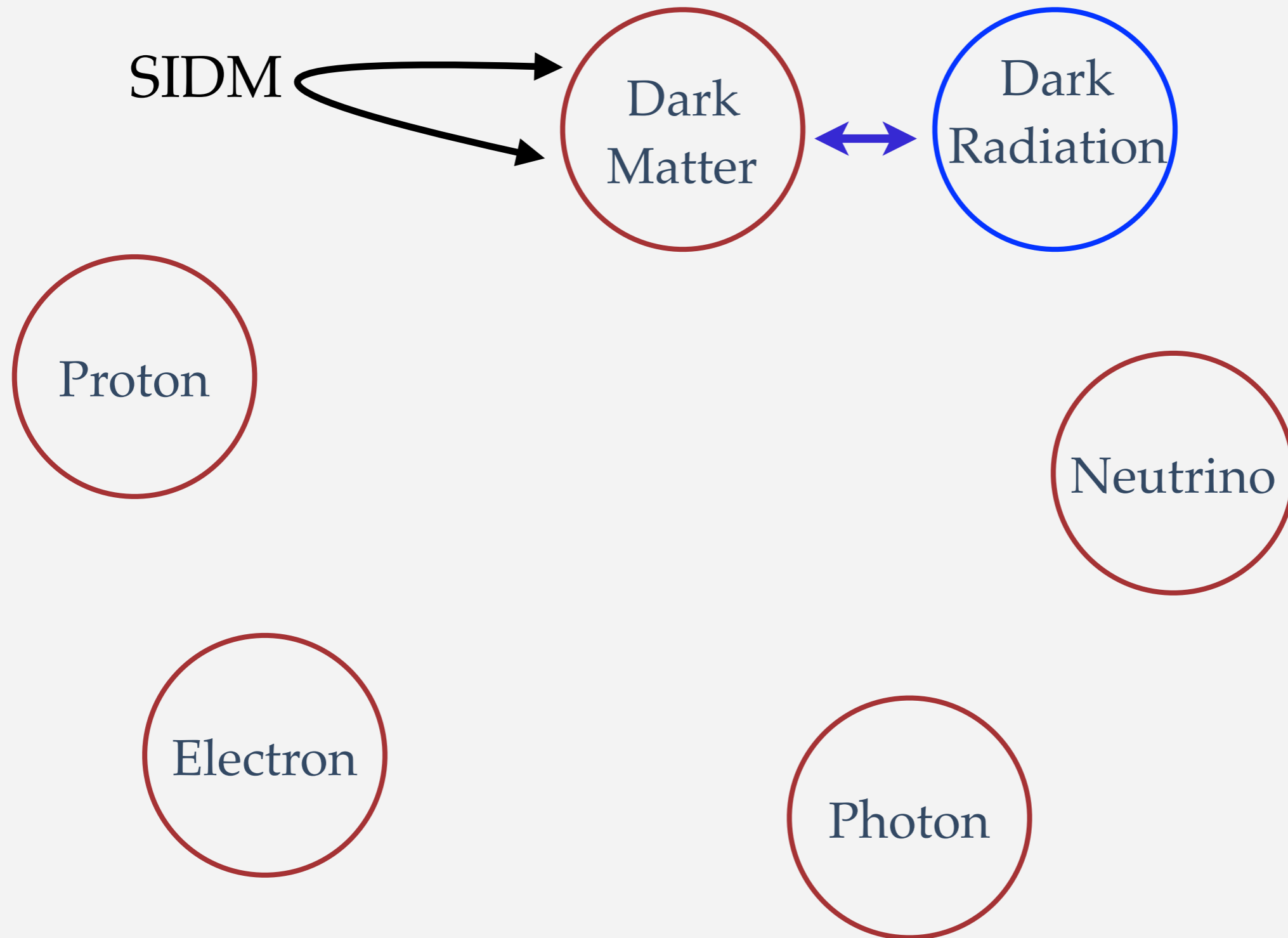
arXiv: 1611.05879, Valentina Prilepina and YT

DM+DE+M/ Anti-M asymmetry, NCTS, Dec 31 2016

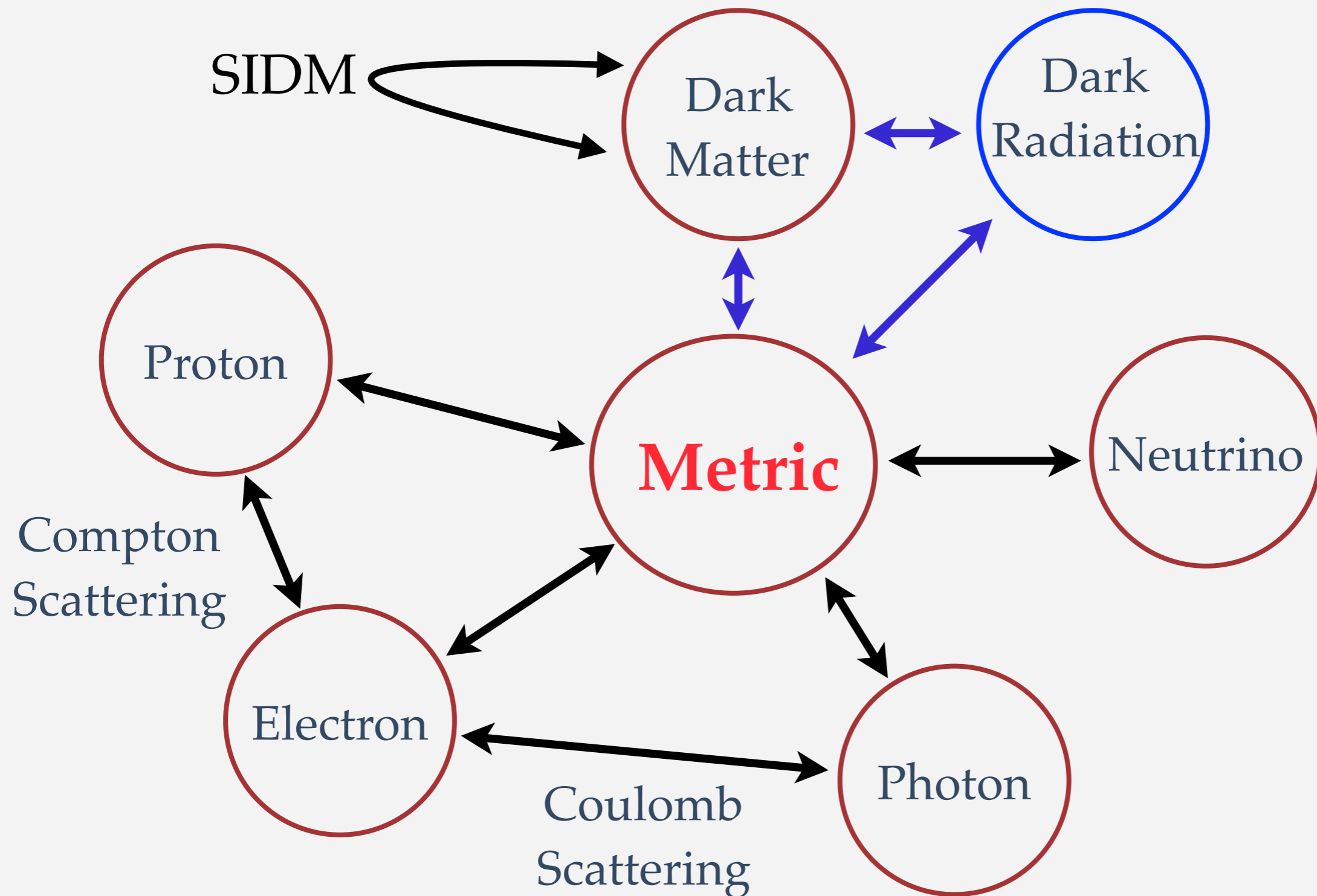
We've been trying very hard to see DM



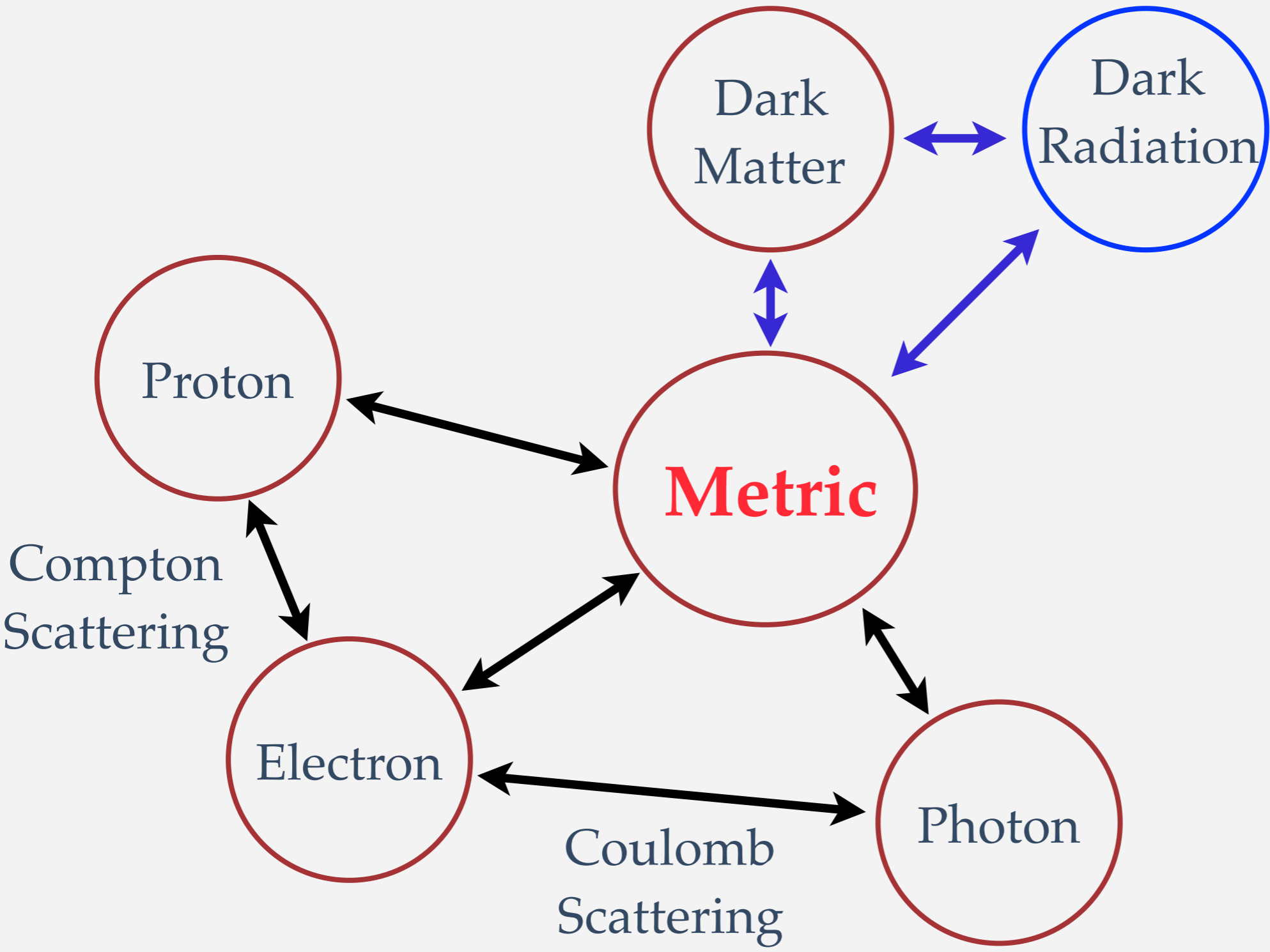
Maybe DM just doesn't couple to SM?



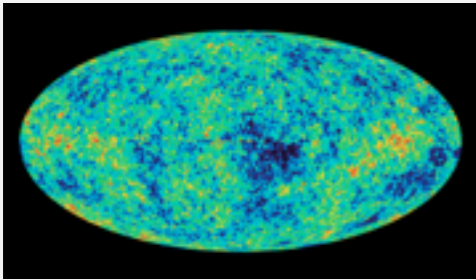
No, everything couples to gravity



Dark sector in Large Scale Structure physics

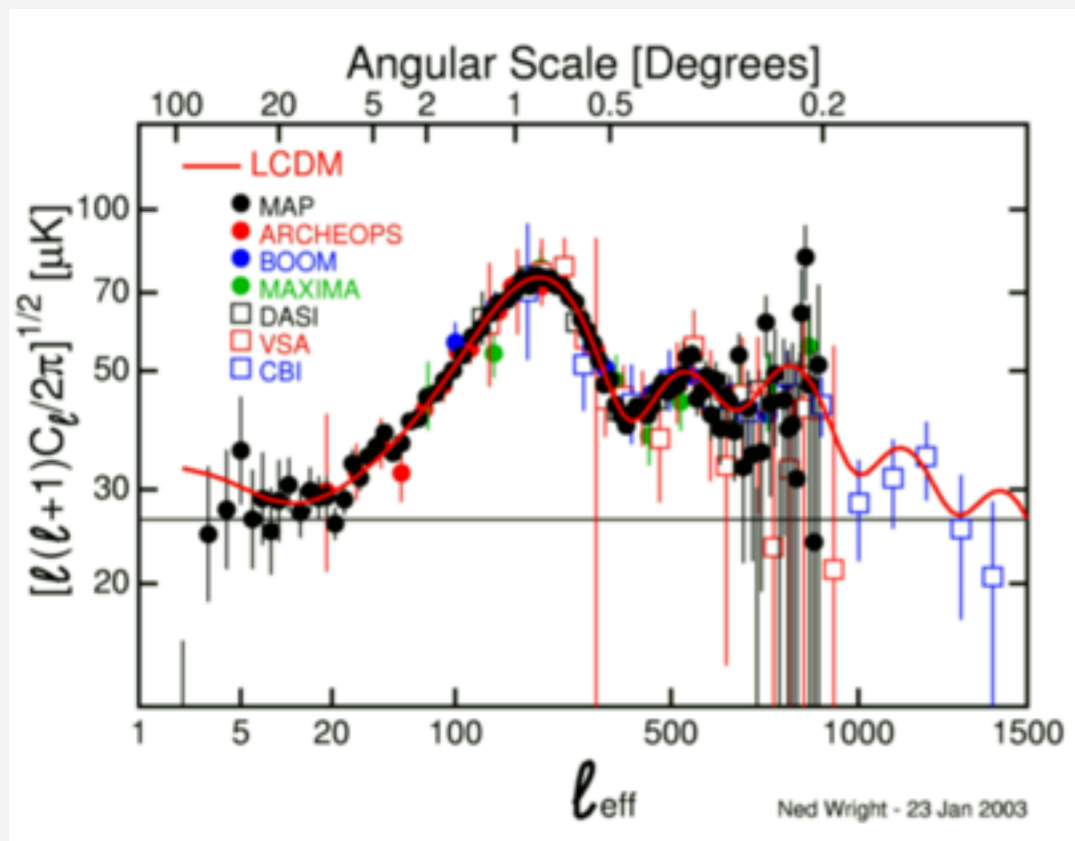


Precision measurement in cosmology

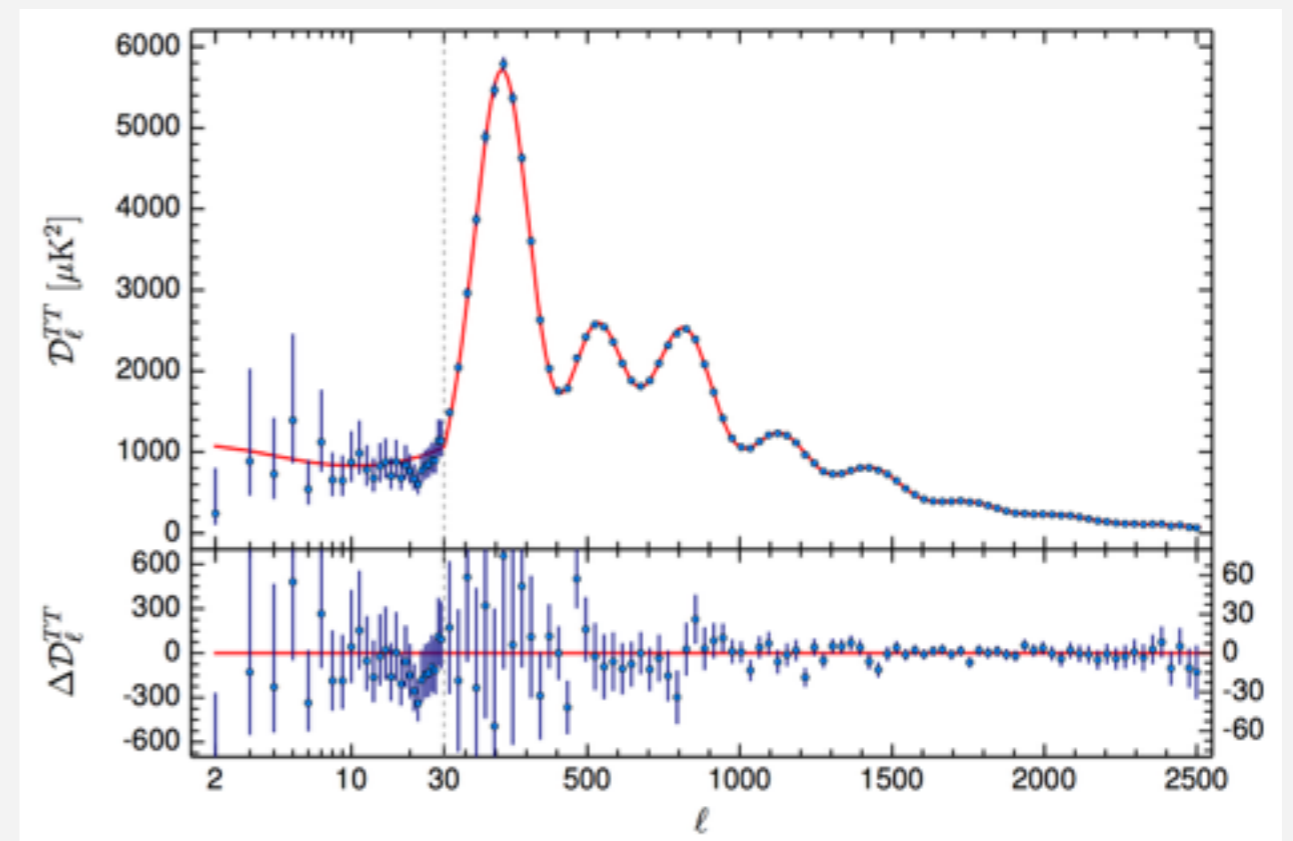


CMB Spectrum

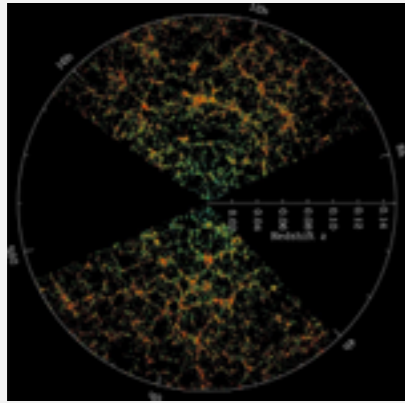
2003



2015



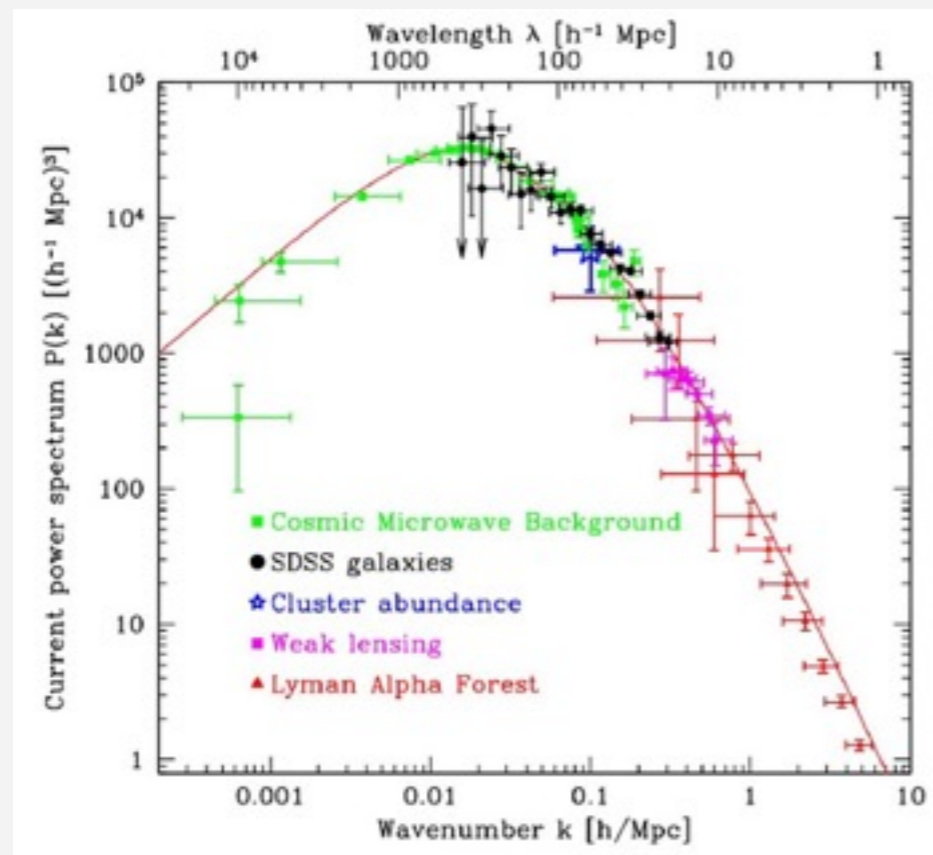
Precision measurement in cosmology



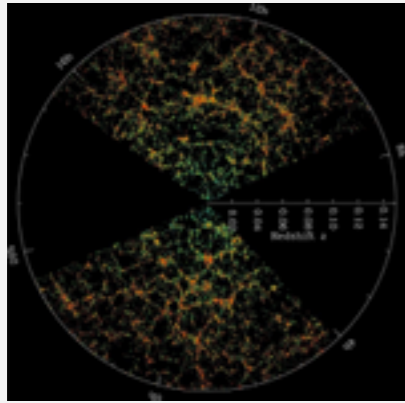
Matter Power Spectrum

Three ways to measure the spectrum

2004



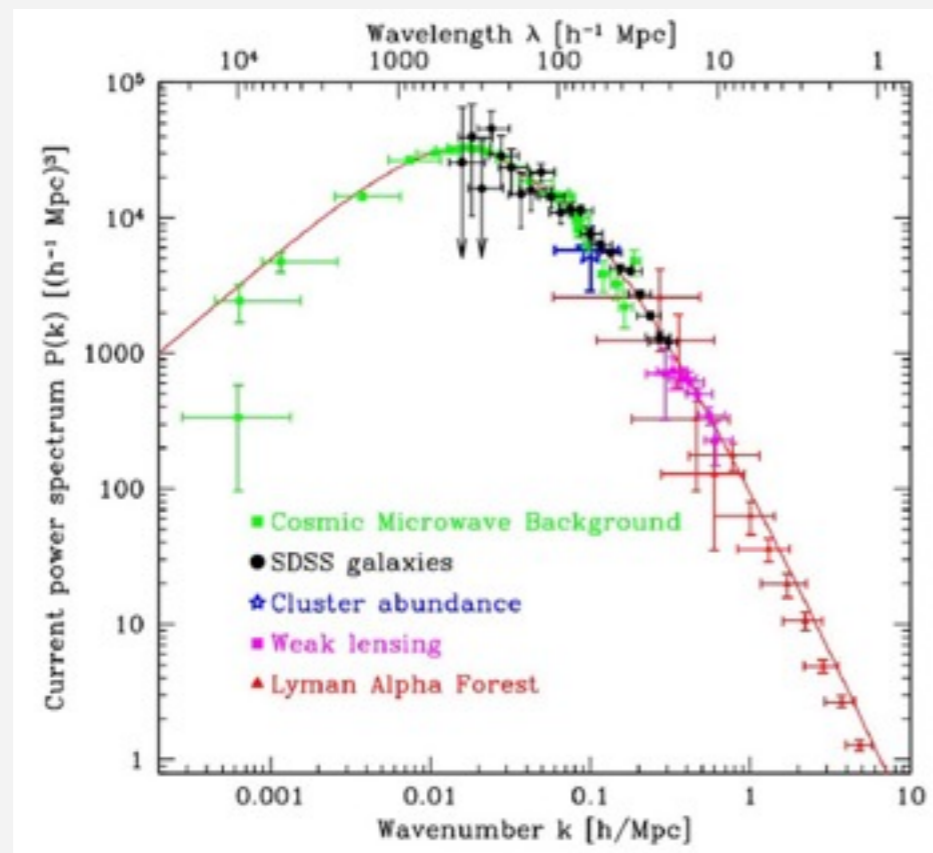
Precision measurement in cosmology



Matter Power Spectrum

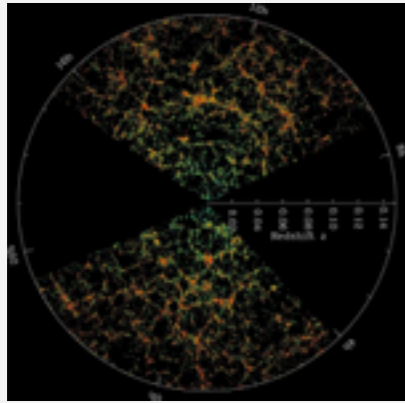
2004

Three ways to measure the spectrum



Assuming a DM model (Λ CDM), fix the parameters using CMB, predict the power spectrum today

Precision measurement in cosmology



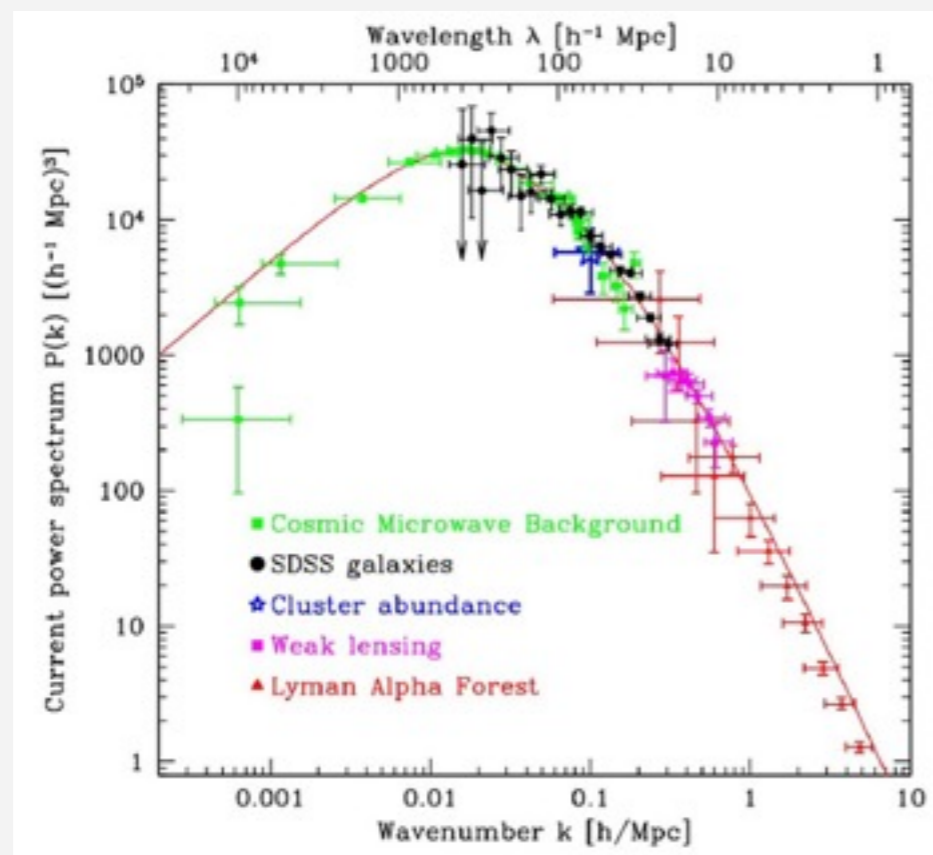
Matter Power Spectrum

2004

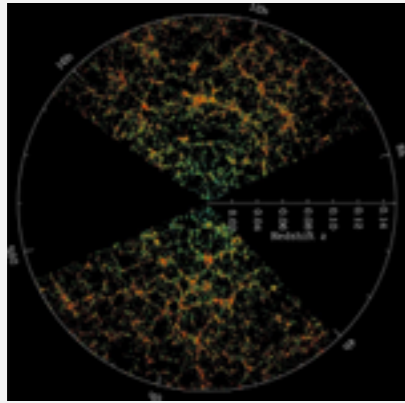
Three ways to measure the spectrum

Assuming a DM model (Λ CDM), fix the parameters using CMB, predict the power spectrum today

Map the galaxy distribution, then fit the DM distribution



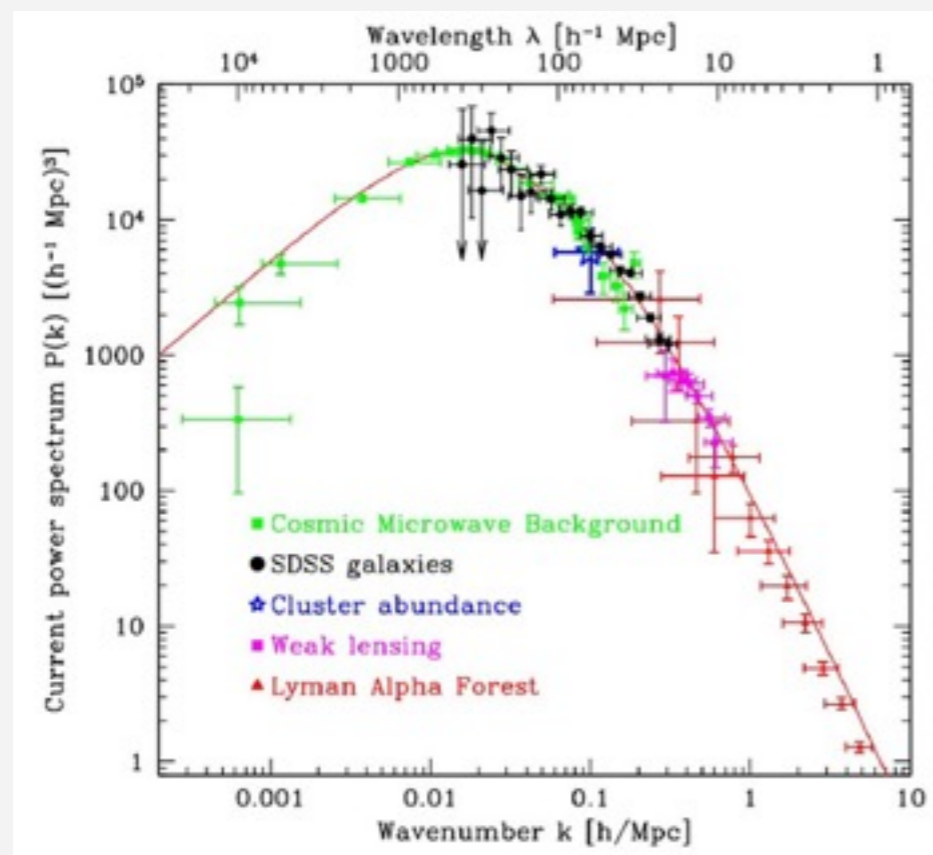
Precision measurement in cosmology



Matter Power Spectrum

2004

Three ways to measure the spectrum

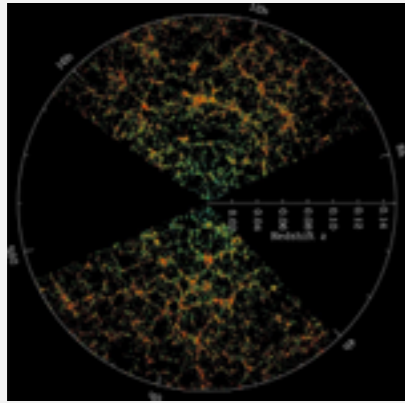


Assuming a DM model (Λ CDM), fix the parameters using CMB, predict the power spectrum today

Map the galaxy distribution, then fit the DM distribution

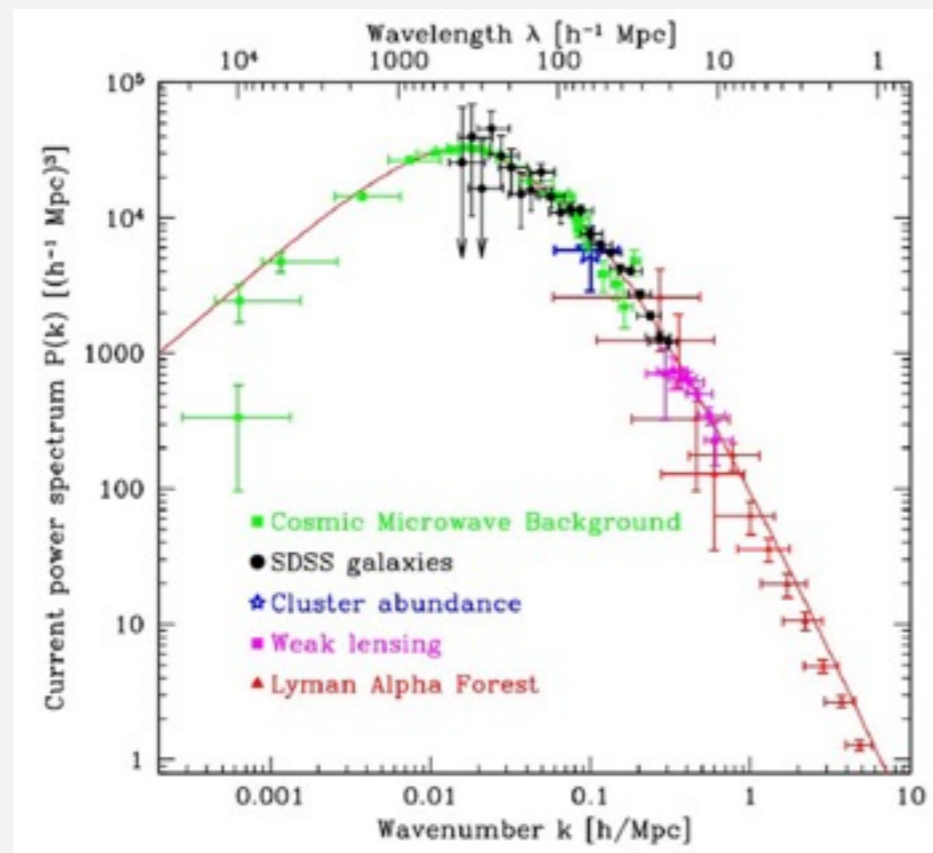
Map the DM distribution directly using weak lensing experiments

Precision measurement in cosmology

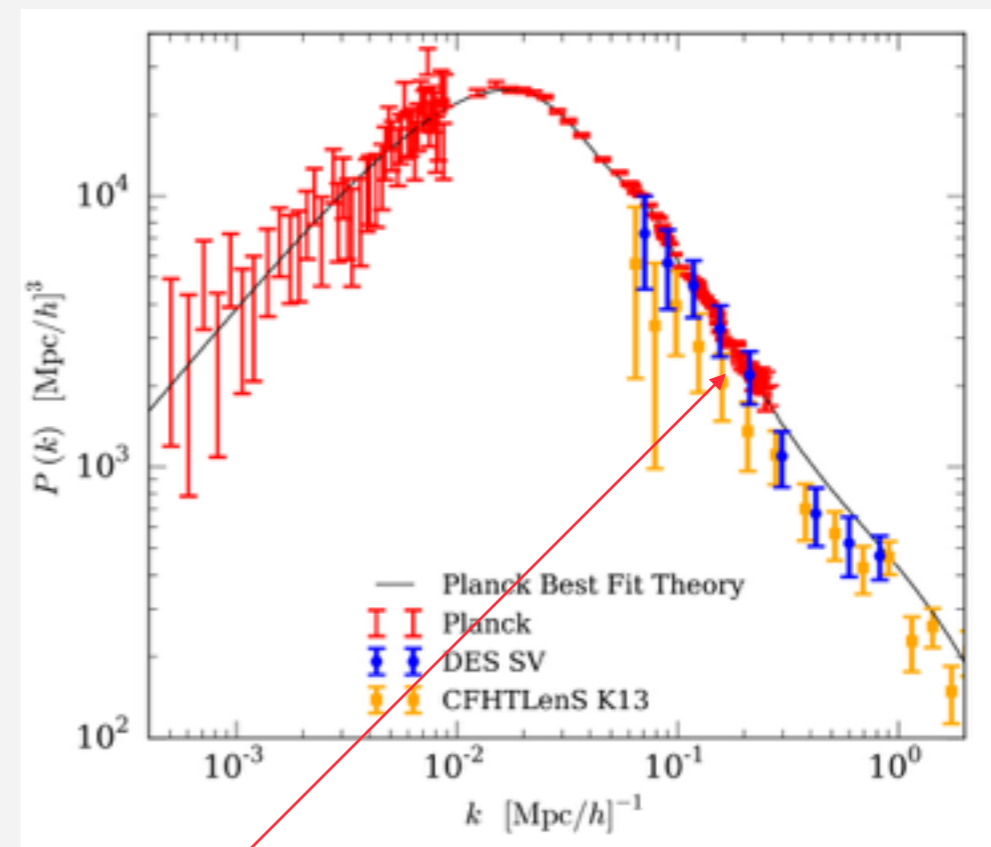


Matter Power Spectrum

2004

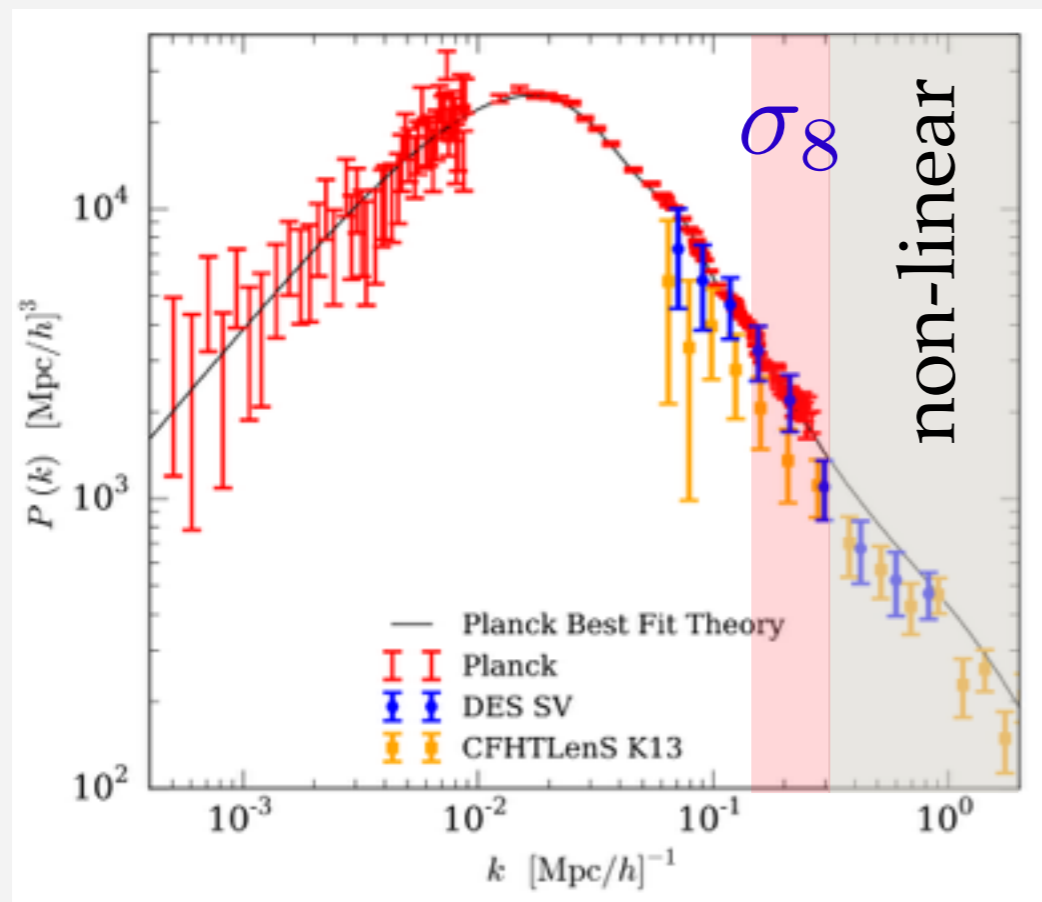


2015



Can start to cross check the CMB & WL results

The σ_8 problem



DES: 1507.05552

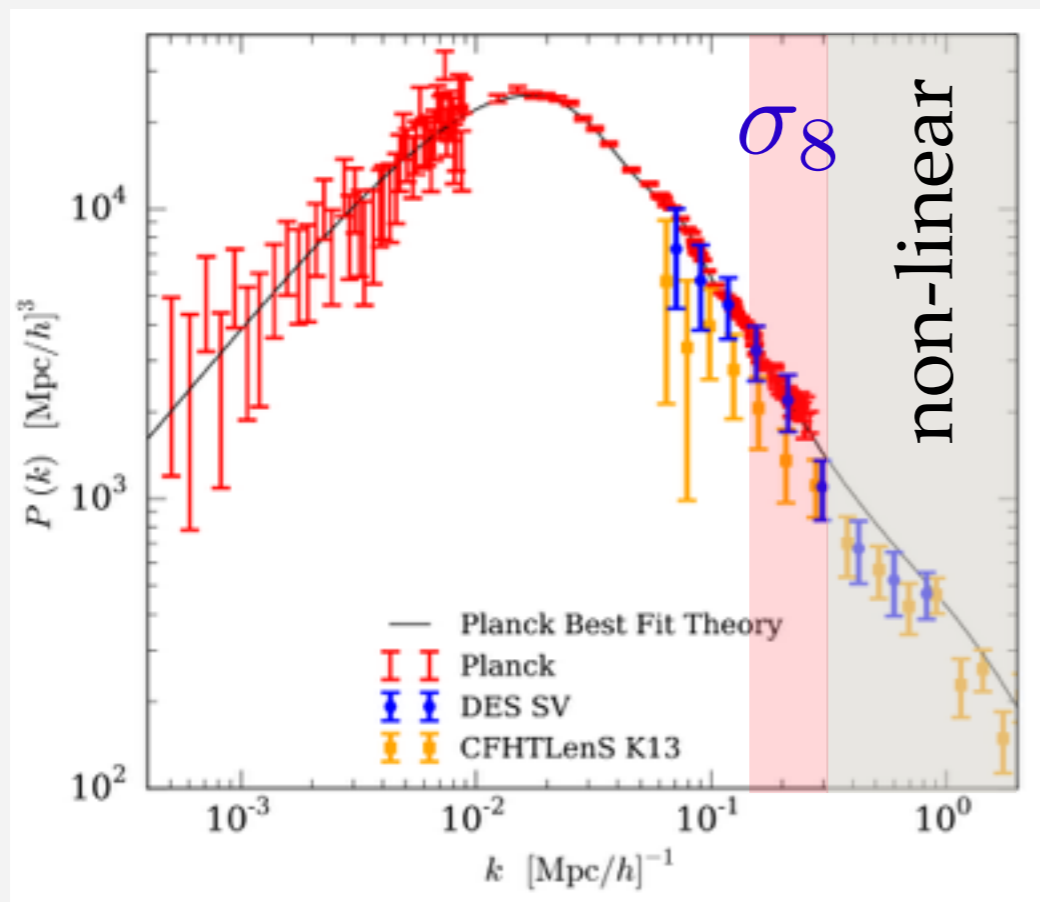
σ_8

~ amplitude of matter fluctuation
on the scale of $8 h^{-1} \text{Mpc}$.

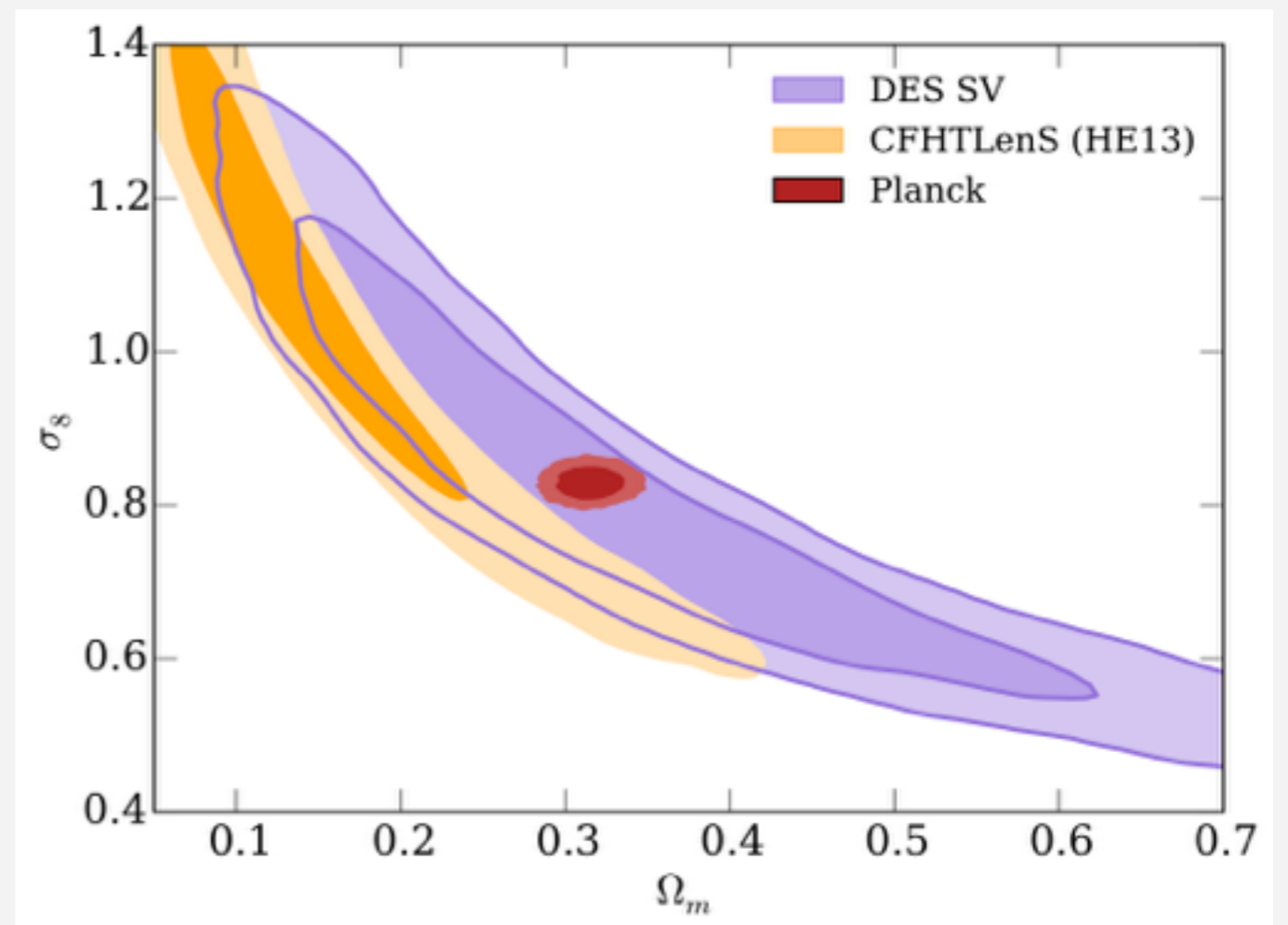
The smallest structure to study without
significant non-linearity effects

The σ_8 problem

Two σ_8 measurements: **CMB + Λ CDM** vs. **Weak Lensing**



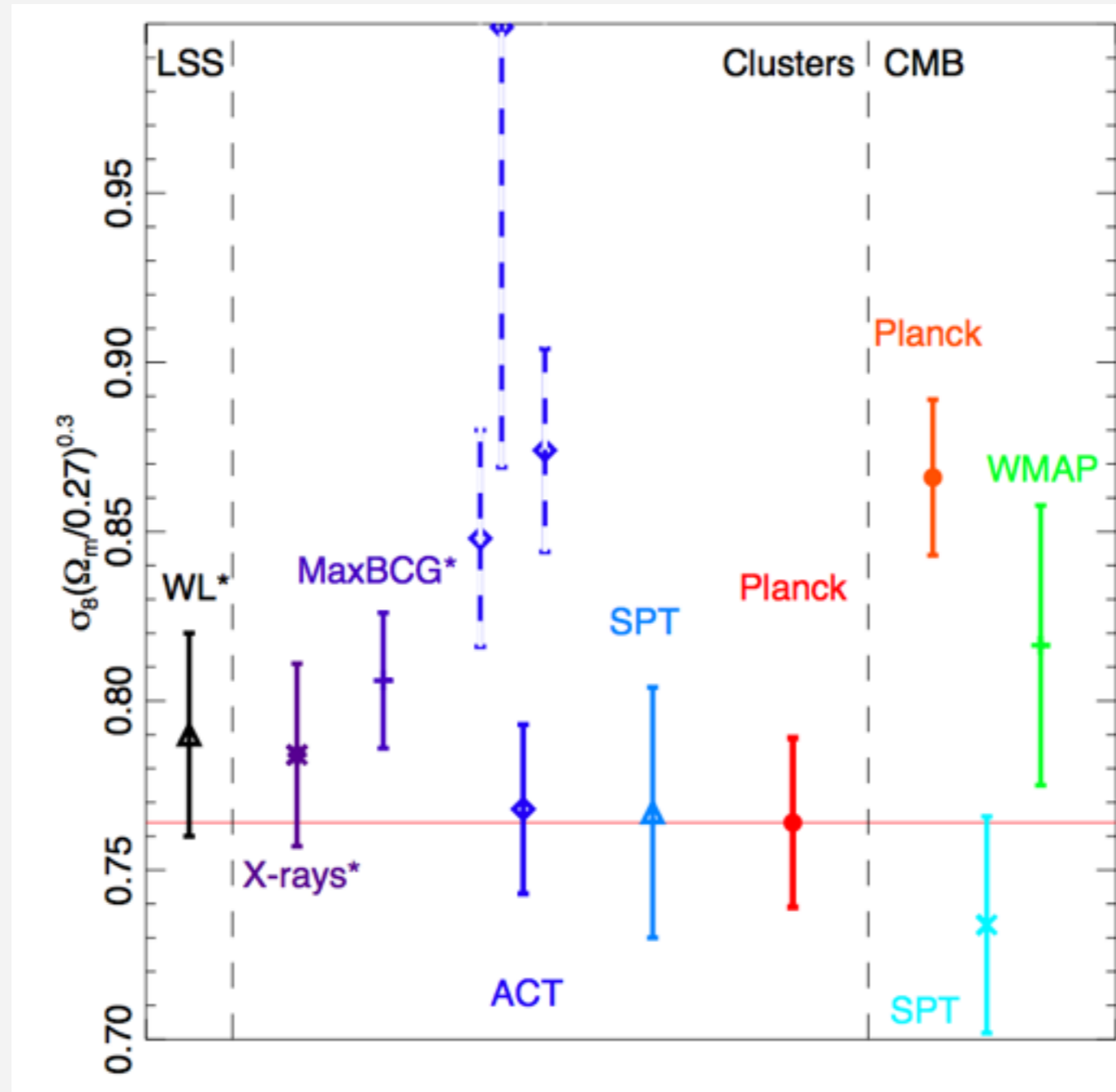
DES: 1507.05552



The CFHTLenS & CMB results deviate by $\sim 2 - 3\sigma$.

Results from galaxy counts

Planck 1303.5080



H₀ problem

Two H₀ measurements

CMB + Λ CDM .

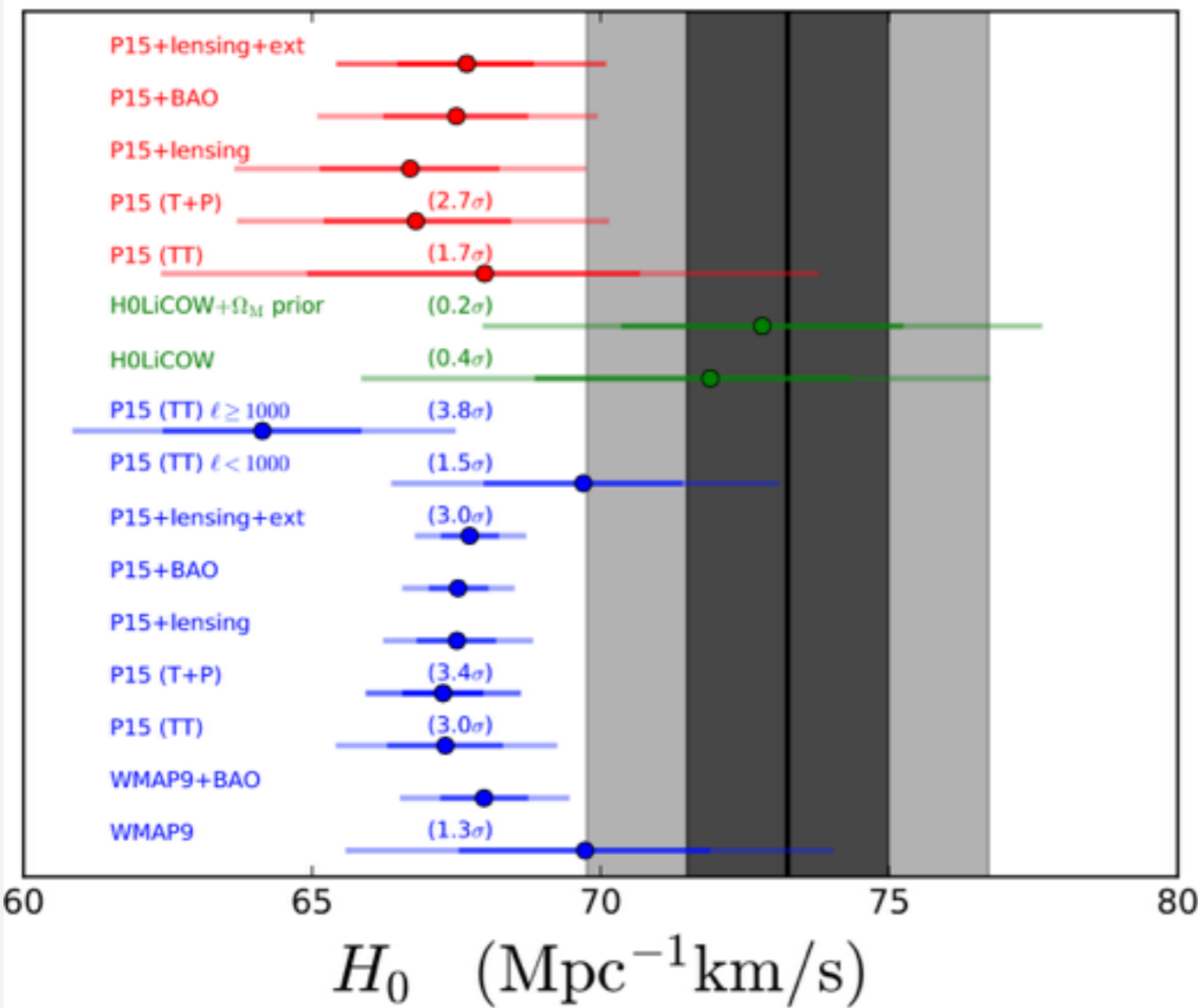
vs.

Local Measurements

$$H_0^{\text{Planck}} = 67.3 \pm 0.7 \text{ km s}^{-1} \text{ Mpc}^{-1}$$

$$H_0^{\text{HST}} = 73.02 \pm 1.79 \text{ km s}^{-1} \text{ Mpc}^{-1}$$

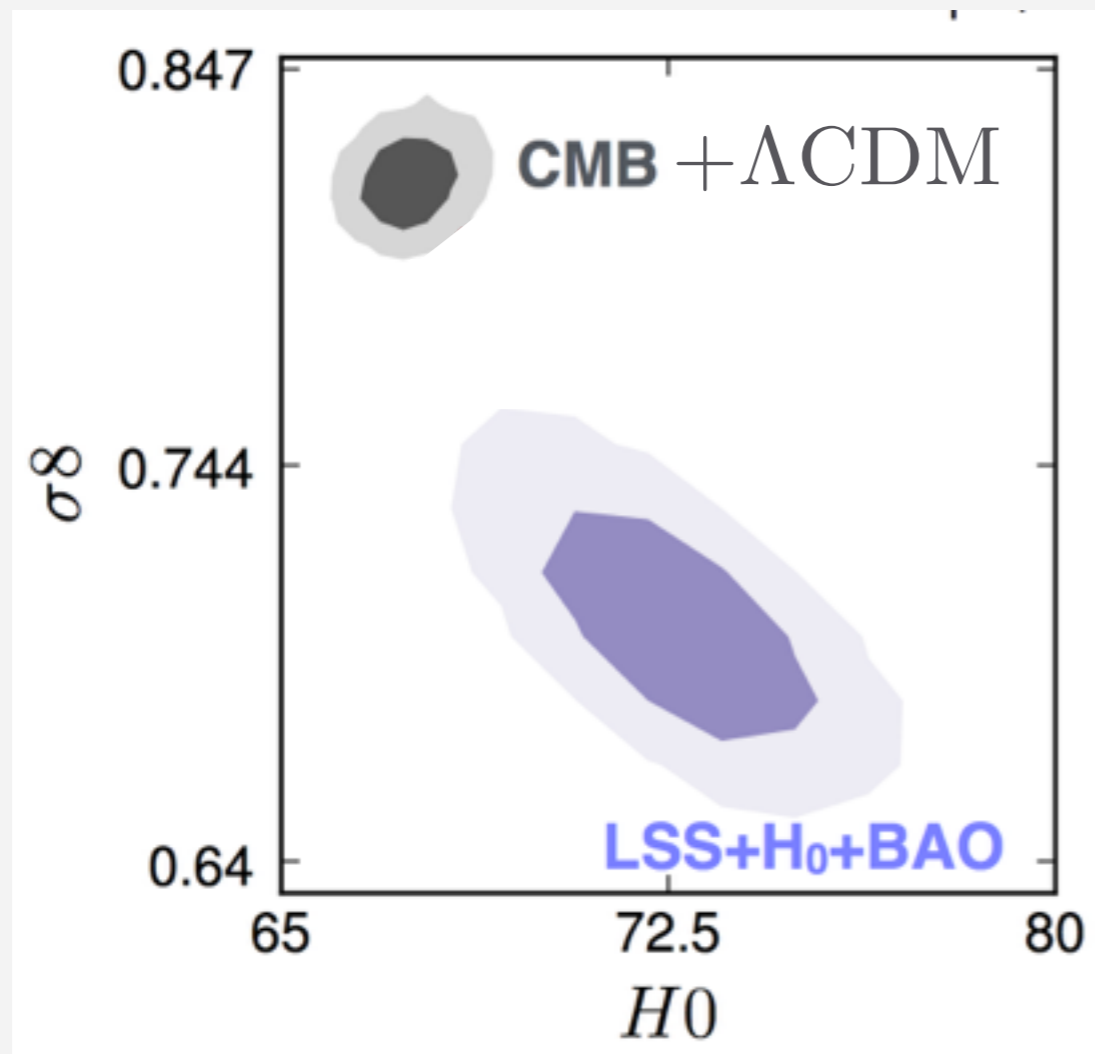
> 3 σ Discrepancy



Bernal et. al. 1607.05617

Puzzles of Large Scale Structure

Poulin et. al. 1606.02073



Comparing to LCDM model,
we want to obtain a

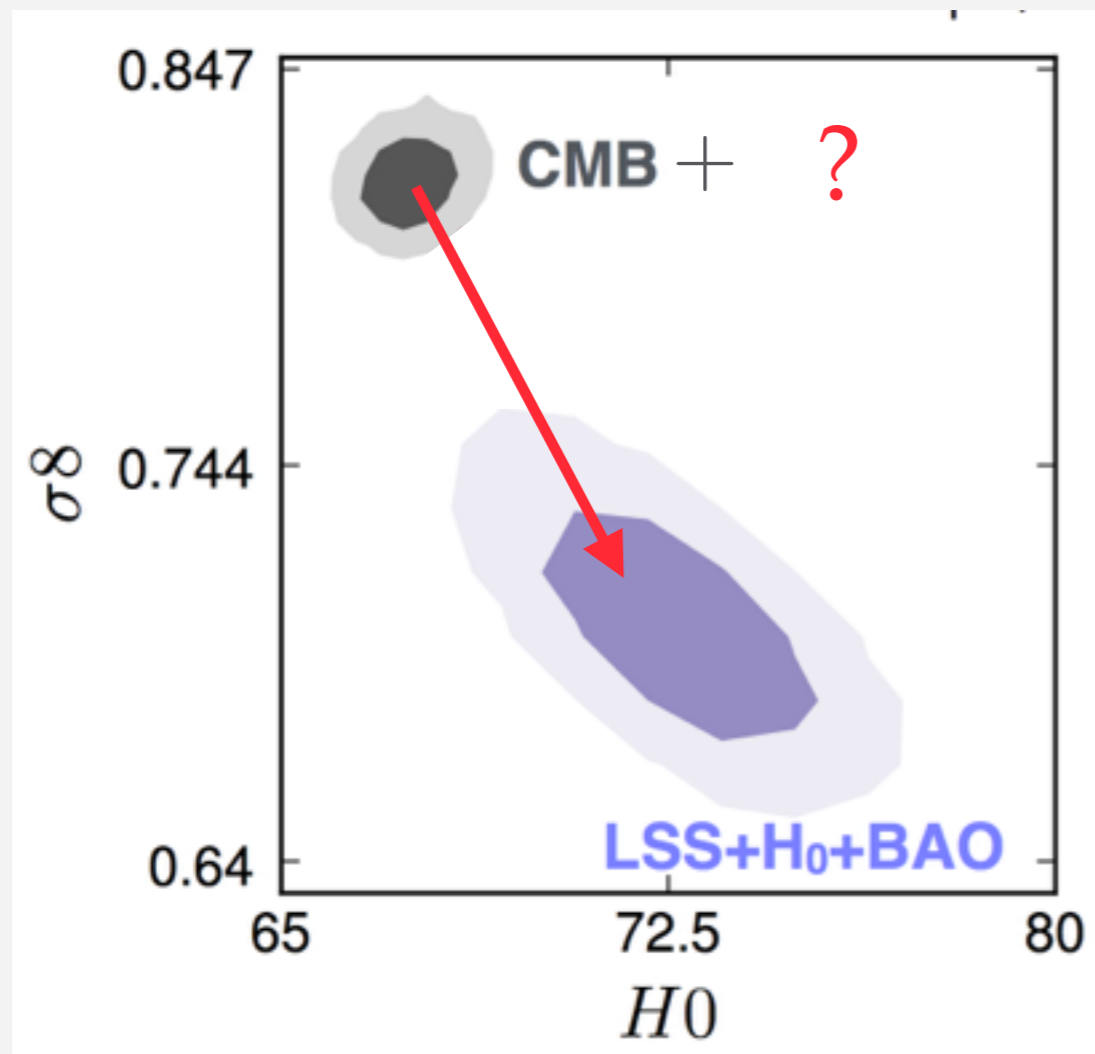
Smaller density perturbation

Larger Hubble expansion

at the late time universe

Puzzles of Large Scale Structure

Poulin et. al. 1606.02073



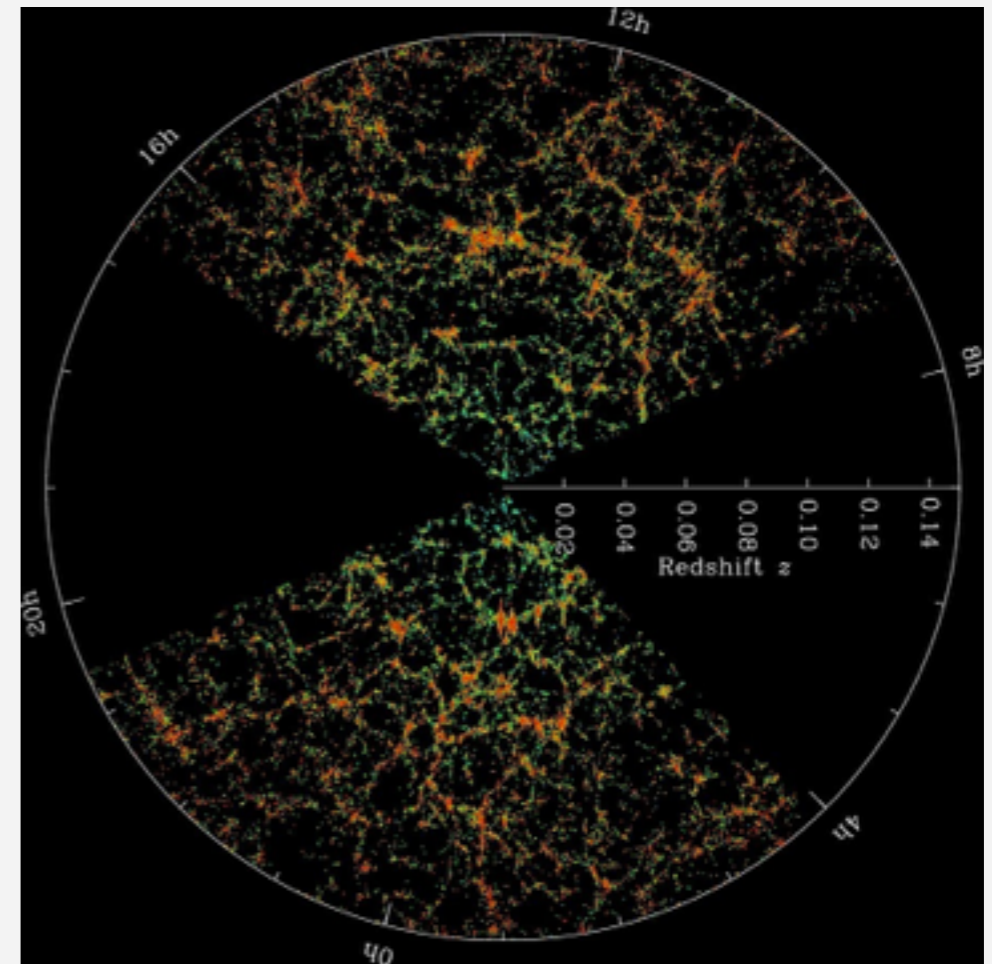
Comparing to LCDM model,
we want to obtain a

Smaller density perturbation

Larger Hubble expansion

at the late time universe

One solution: Partially Acoustic DM



DISTRIBUTION OF GALAXIES IN OUR UNIVERSE. CREDIT: SDSS

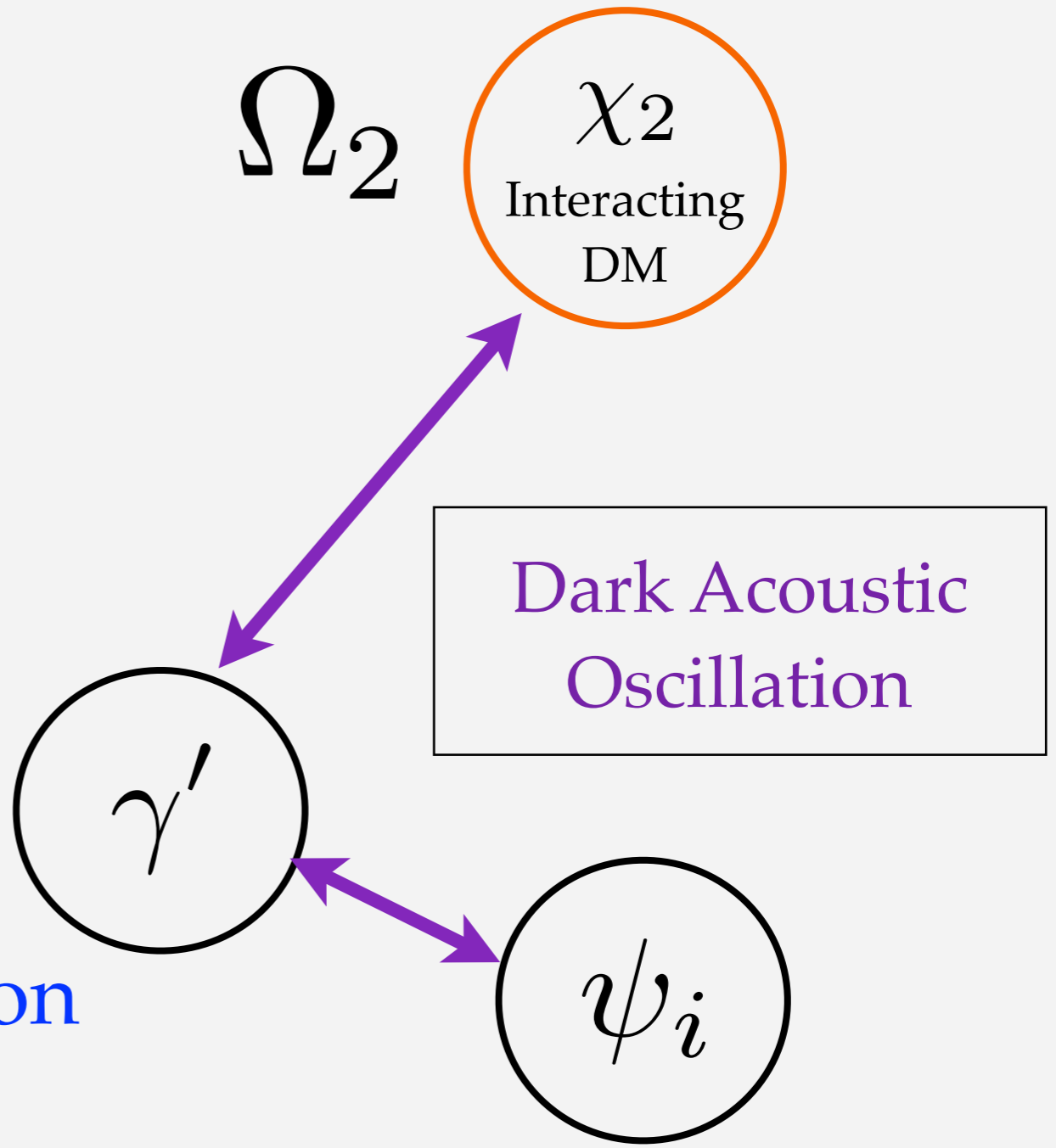
Consider a non-minimal dark sector



$$r \equiv \Omega_2 / \Omega_{\text{DM}}$$

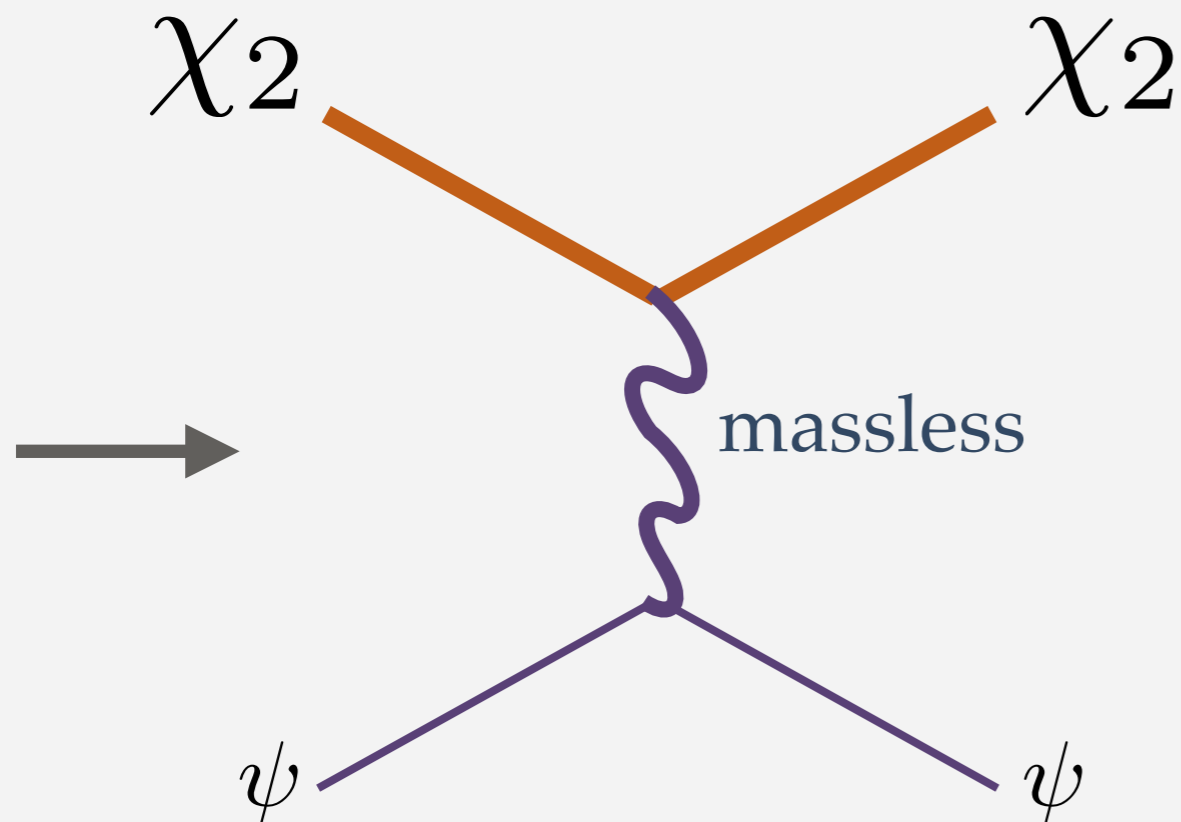
$$r < 1$$

Partially Acoustic Oscillation
(PAcDM)



For the acoustic oscillation to exist

We need the DM-DR scattering to remain non-decoupled



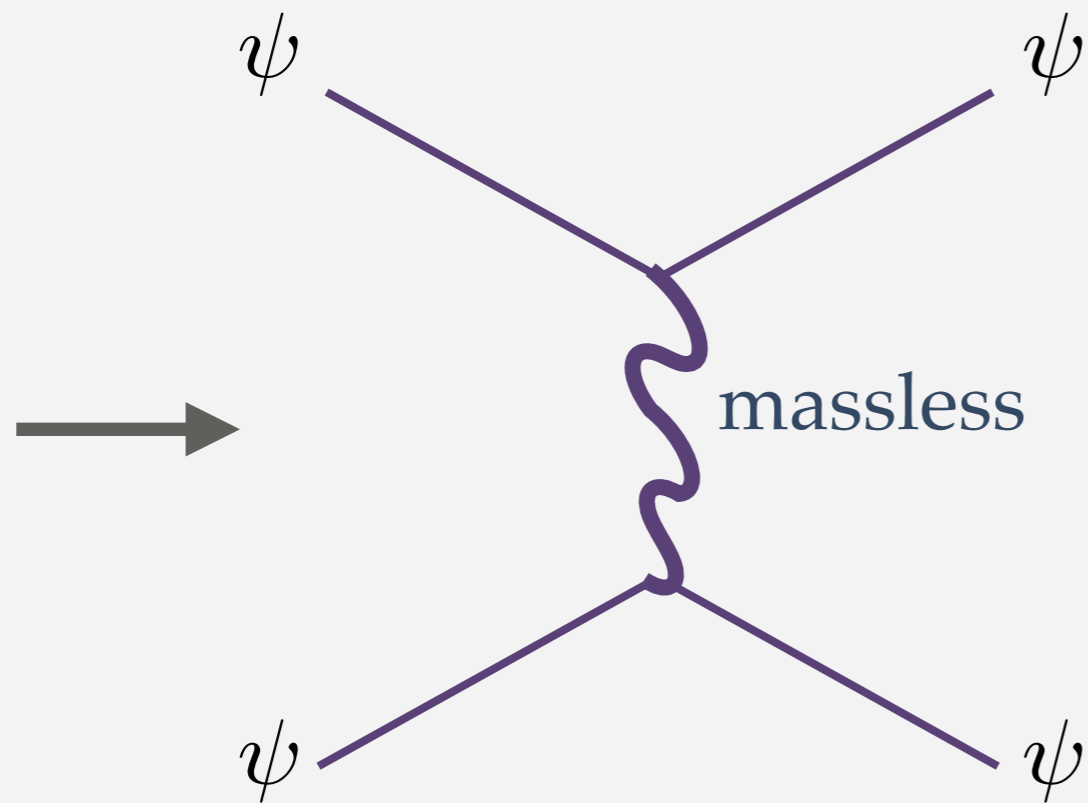
$$\Gamma \simeq \hat{\alpha}^2 \ln(\hat{\alpha}^{-1}) \frac{T_D^2}{m_{\text{DM}}}$$

Same temp-dependence as Hubble
in the radiation-dominant era

Easy to keep $\Gamma \gg H$ all the time, if $\hat{\alpha} \gg 10^{-8} \sqrt{\frac{m_{\chi_2}}{10 \text{ GeV}}}$

Tightly coupled dark radiation

We need the DM-DR scattering to remain non-decoupled

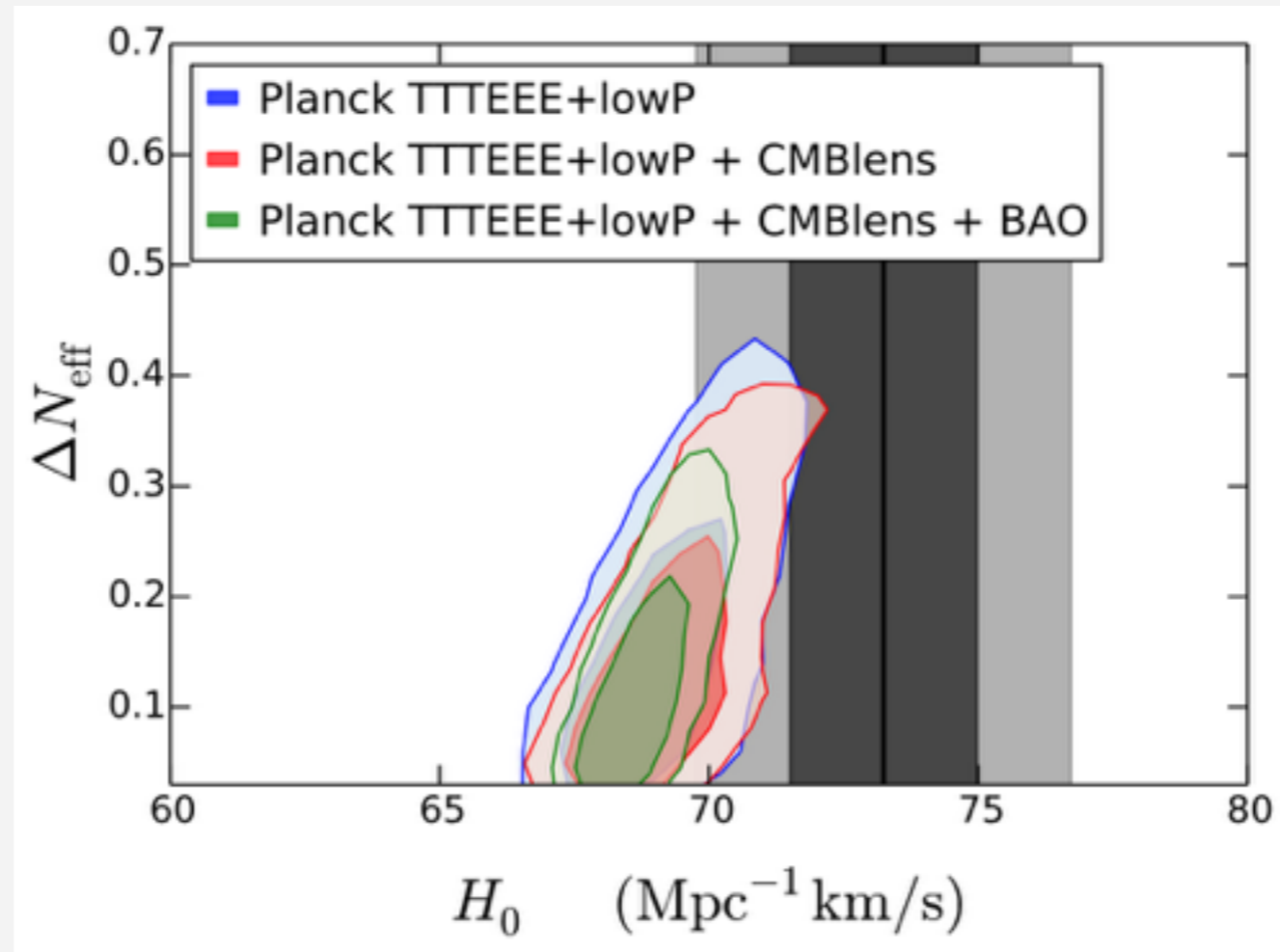


$$\Gamma \simeq \hat{\alpha}^2 \ln(\hat{\alpha}^{-1}) T_D$$

The same coupling keeps dark fermions / photon a tightly coupled fluid

Solving H_0 problem with extra dark radiation

Bernal et. al. 1607.05617

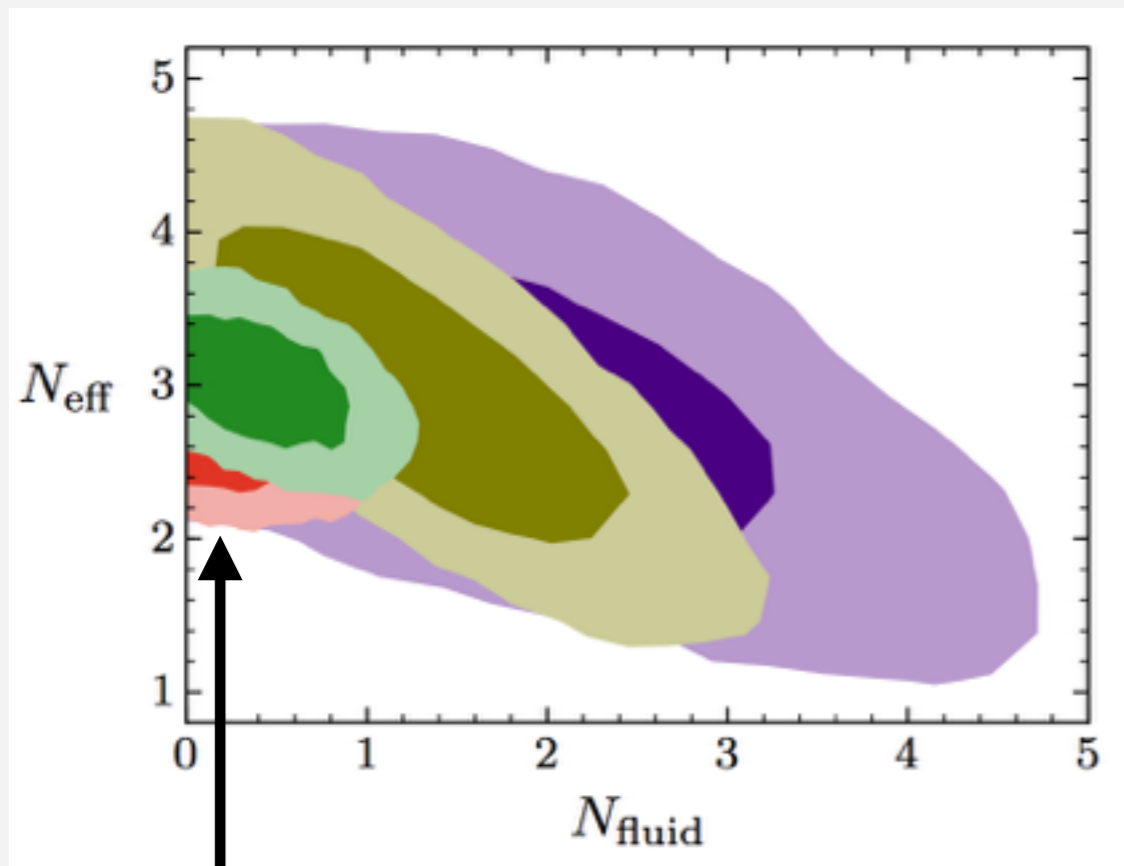


Can explain the larger H_0 by including $\Delta N_{\text{eff}} > 0.4$ dark radiation

Adam Riess et.al. 1604.01424

Dark fluid is better than FS-radiation

Baumann et. al. 1508.06342



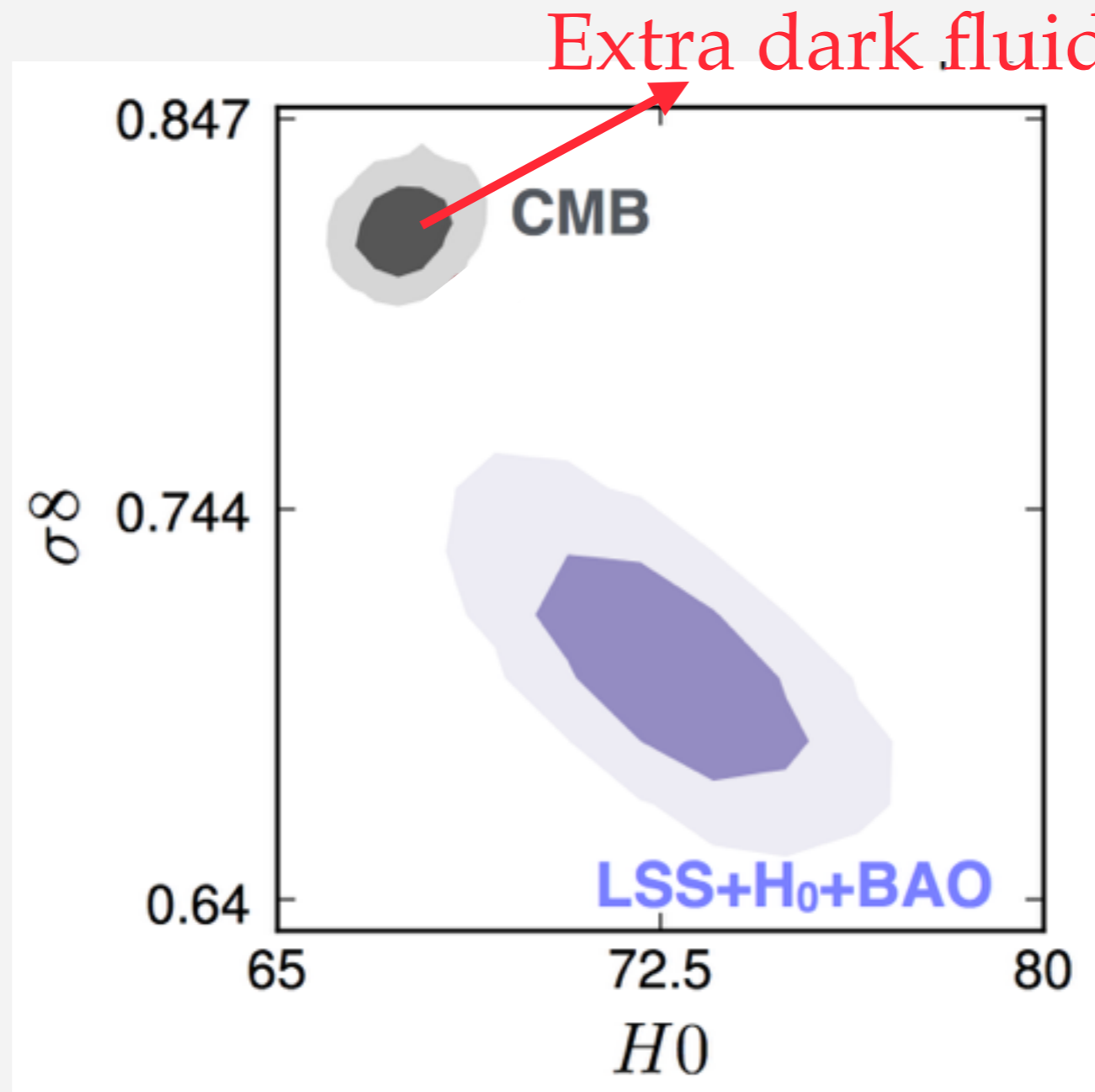
Planck TT, TE, and EE likelihoods

ΔN_{eff} bound on a tightly coupled fluid is weaker

	TT, TE, EE		TT-only	
	varying Y_p	fixed Y_p	varying Y_p	fixed Y_p
N_{eff}	$2.78^{+0.30}_{-0.35}$	$2.99^{+0.30}_{-0.29}$	$2.87^{+0.76}_{-0.74}$	$2.94^{+0.71}_{-0.69}$
N_{fluid}	< 0.88	< 1.06	< 3.93	< 2.65

(2σ)

Reconcile H_0 , but makes σ_8 worse



Increase radiation
density

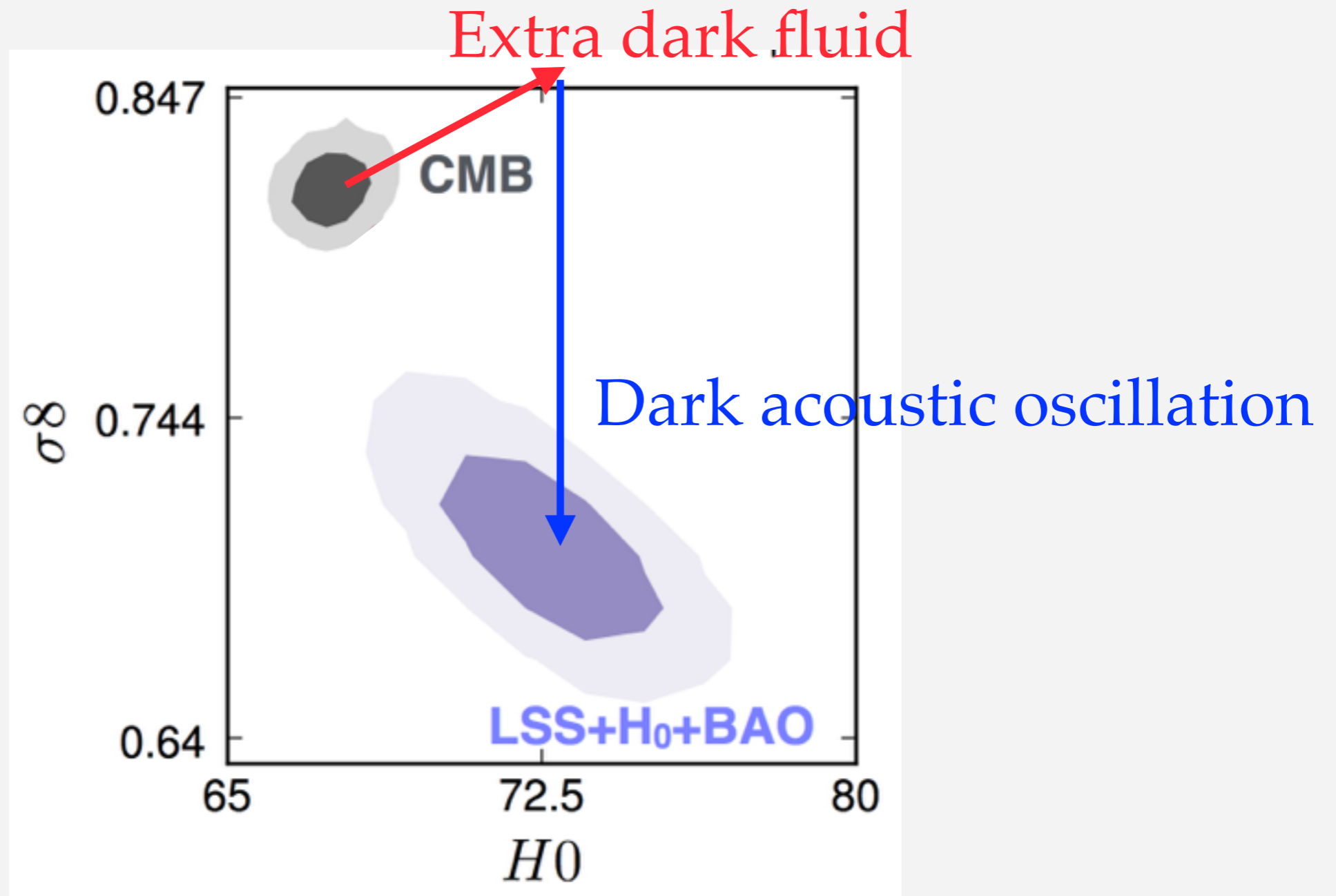


Increase matter
density

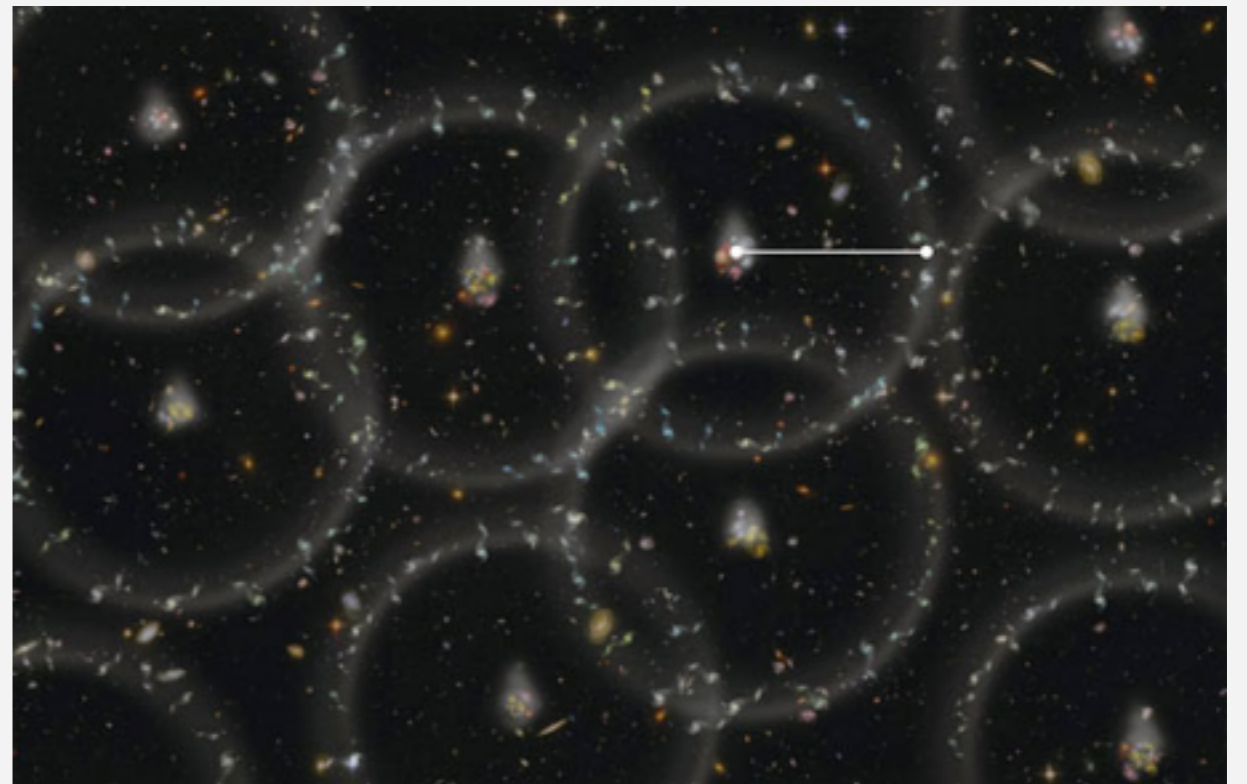


Increase density
perturbation

DM-DR scattering suppresses Σ_8

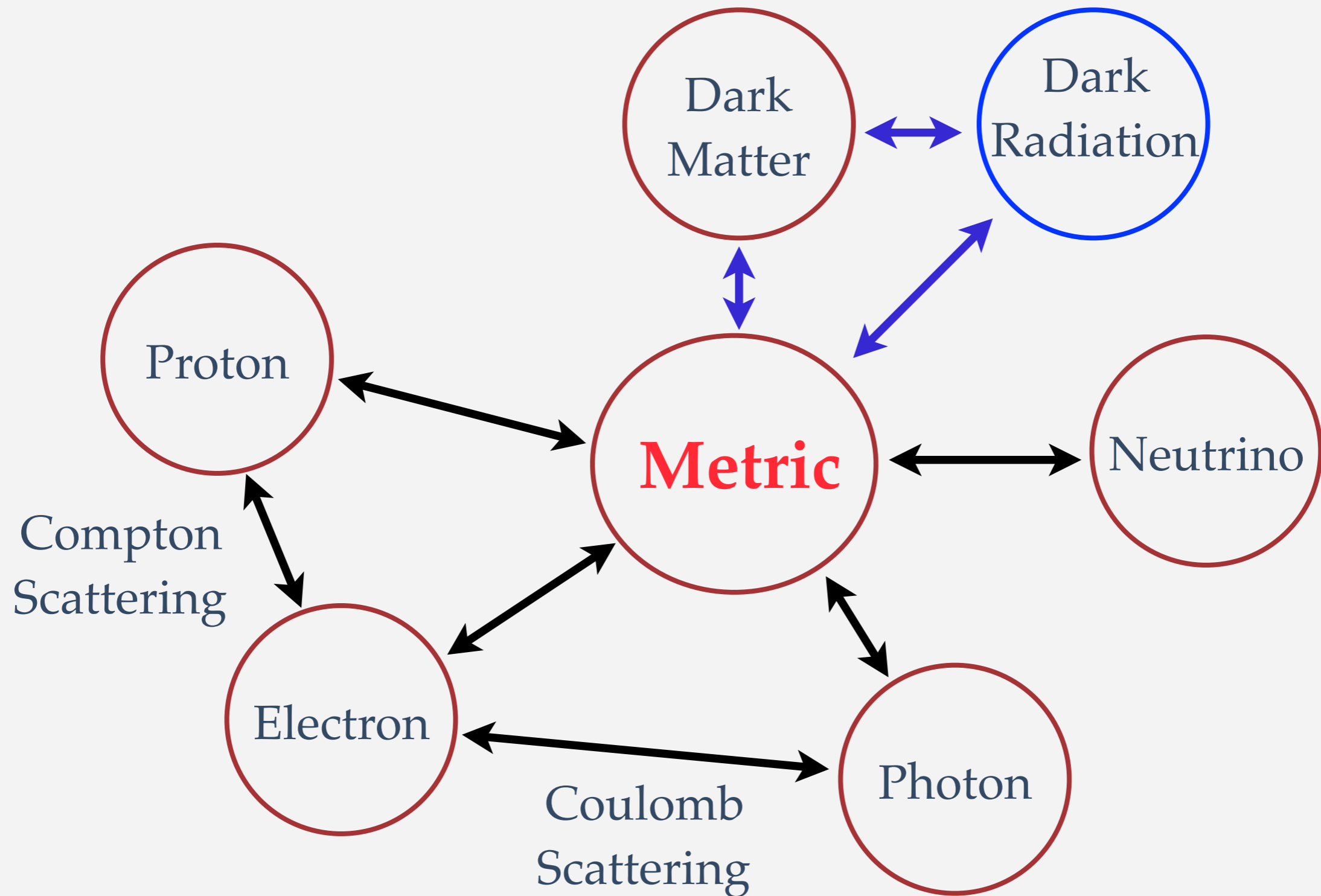


Structure Formation with DAO



A cartoon produced by the BOSS project showing the spheres of baryons around the initial dark matter clumps

Evolution of the Large Scale Structures



In the tightly coupled DM-DR limit

We can simplify the evolution of DM perturbation

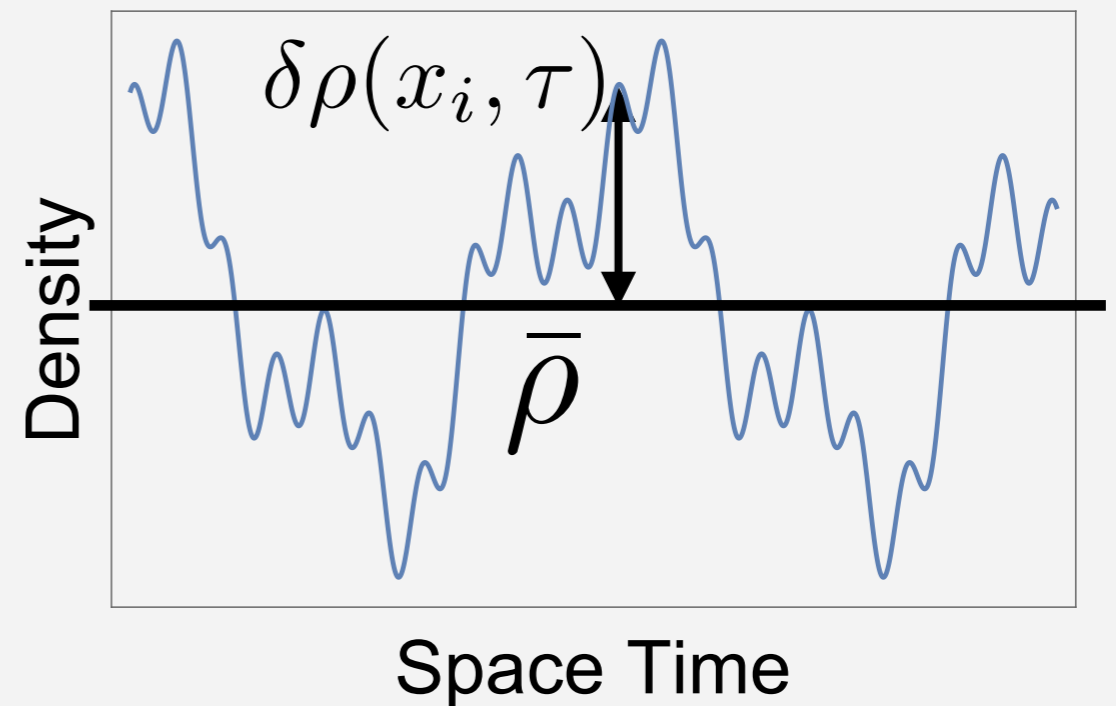
$$\ddot{\delta}_D + \frac{\dot{a}}{a} \frac{R}{1+R} \dot{\delta}_D + \frac{k^2}{3(1+R)} \delta_D \simeq -k^2 \psi$$

metric perturbation

Density Perturbation

$$\delta_D \equiv \frac{\delta \rho_D}{\bar{\rho}_D}$$

$$P(k)_s \propto k^{-3} \langle \delta_s(k, a)^2 \rangle$$



In the tightly coupled DM-DR limit

We can simplify the evolution of DM perturbation

$$\ddot{\delta}_D + \frac{\dot{a}}{a} \frac{R}{1+R} \dot{\delta}_D + \frac{k^2}{3(1+R)} \delta_D \simeq -k^2 \psi$$

$$\delta_D \equiv \frac{\delta\rho_D}{\bar{\rho}_D}$$

$$R \equiv \frac{3\rho_D}{4\rho_R}$$

Parametrize the “mass”
of DM-DR fluid

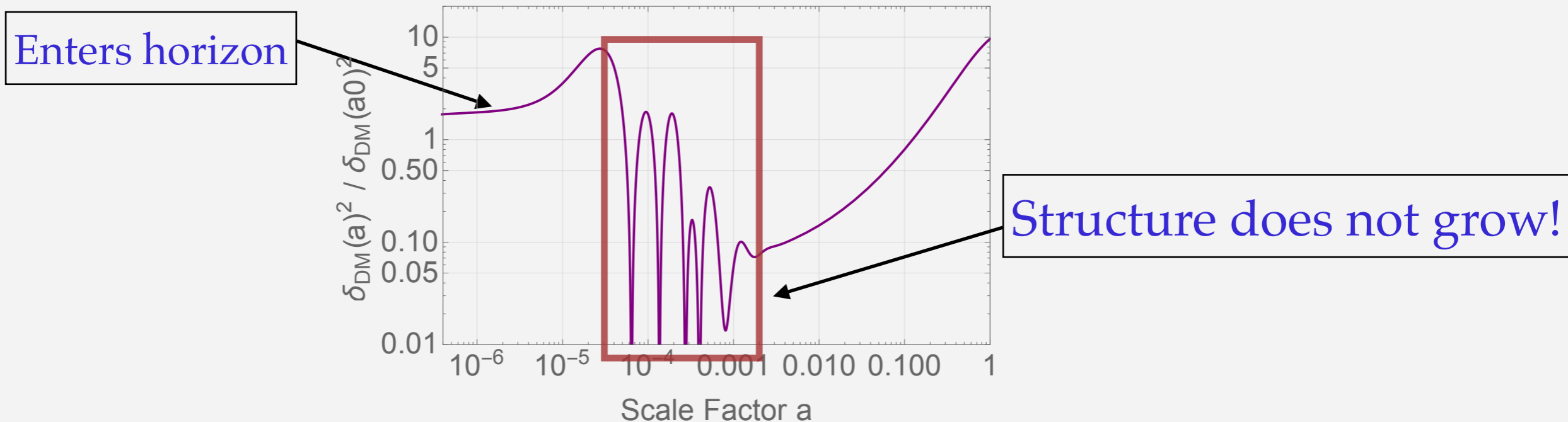
$$P(k)_s \propto k^{-3} \langle \delta_s(k, a)^2 \rangle$$

Radiation Domination , $R \ll 1$

Density perturbation oscillates => No structure grows

$$\ddot{\delta}_D + \frac{\dot{a}}{a} \frac{R}{1+R} \dot{\delta}_D + \frac{k^2}{3(1+R)} \delta_D \simeq -k^2 \psi \quad \text{small in RD}$$

The density perturbation oscillates as a harmonic oscillator!
Same physics as the baryon acoustic oscillation

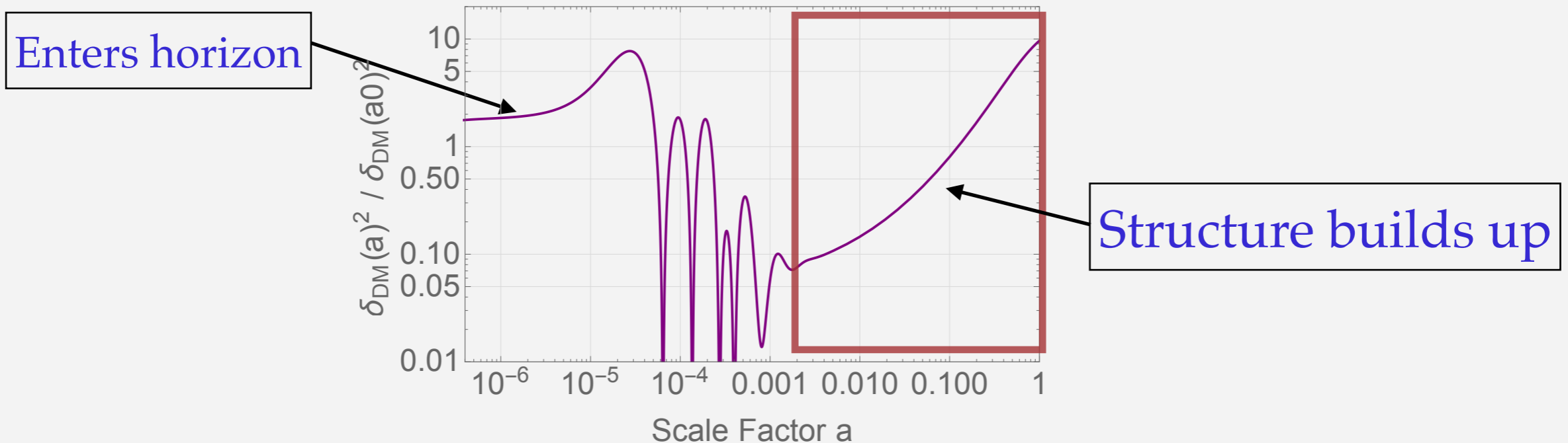


Matter Domination , $R \gg 1$

No oscillation => Linear growth

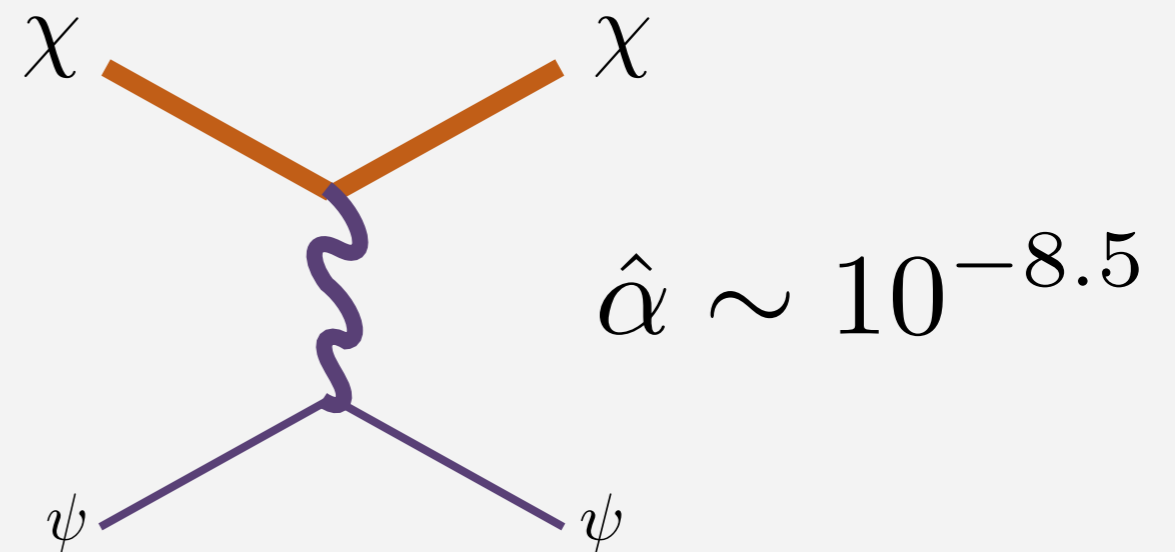
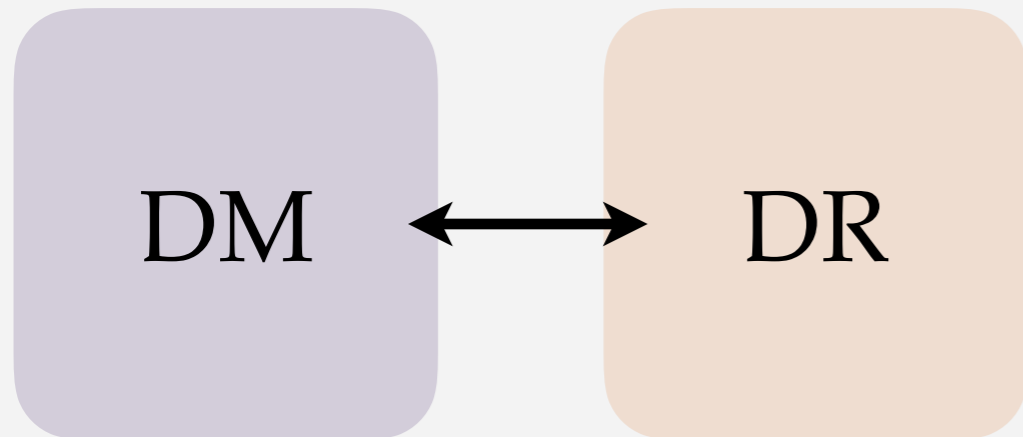
$$\ddot{\delta}_D + \frac{\dot{a}}{a} \frac{R}{1+R} \dot{\delta}_D + \frac{k^2}{3(1+R)} \delta_D \simeq -k^2 \psi$$

No oscillation, no damping from the DR scattering
Same structure formation as cold DM



Quasi-Acoustic Dark Matter

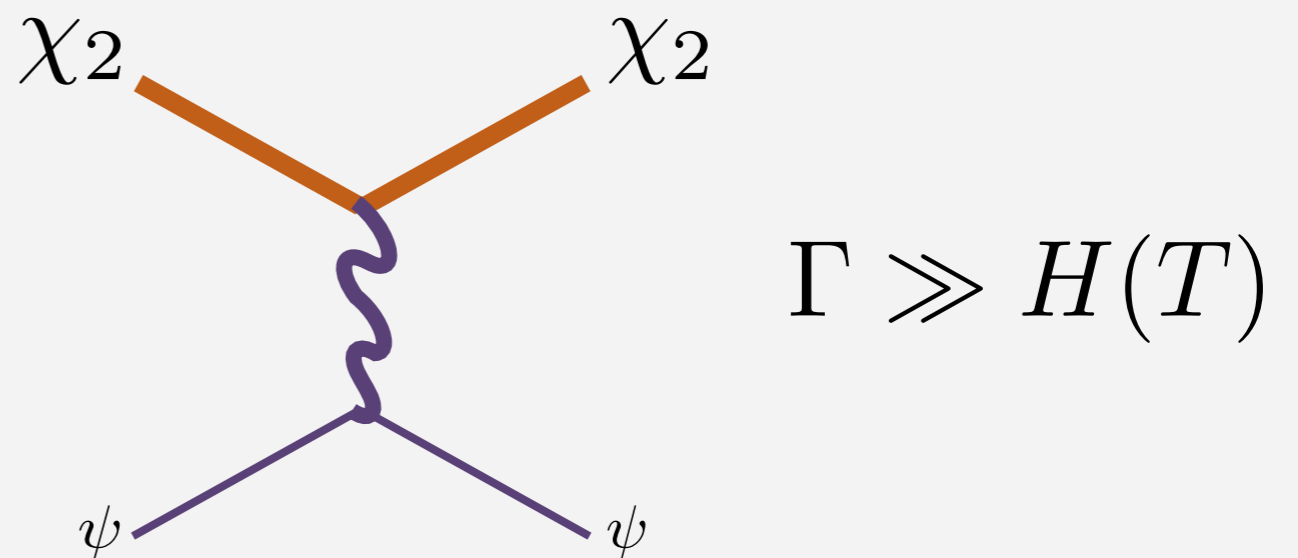
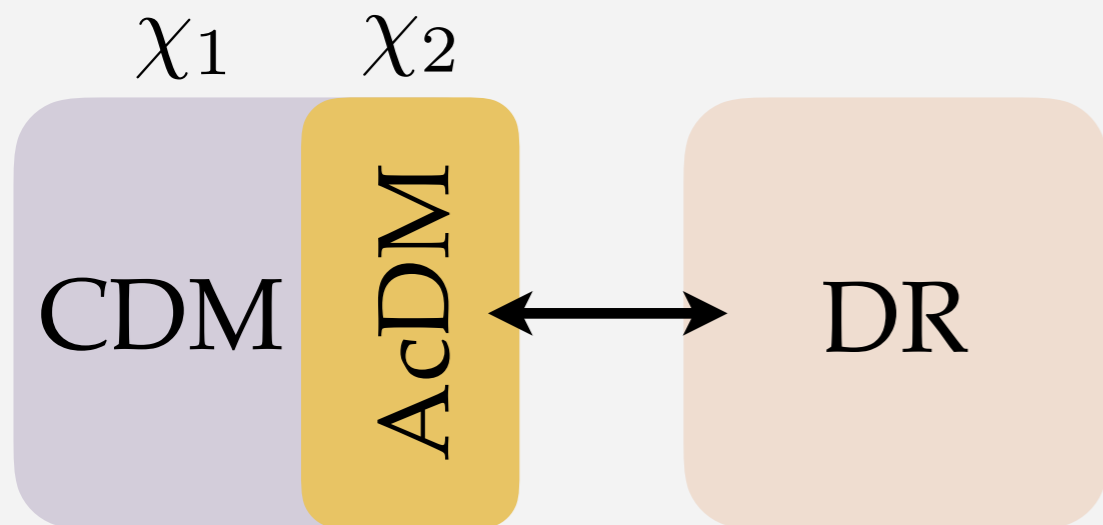
We need a small DM coupling for the right σ_8 suppression



Manuel A. Buen-Abad, Gustavo Marques-Tavares, and Martin Schmaltz (2015)

Julien Lesgourgues, Gustavo Marques-Tavares, and Martin Schmaltz (2015)

Partially-Acoustic Dark Matter



$$\Gamma \gg H(T)$$

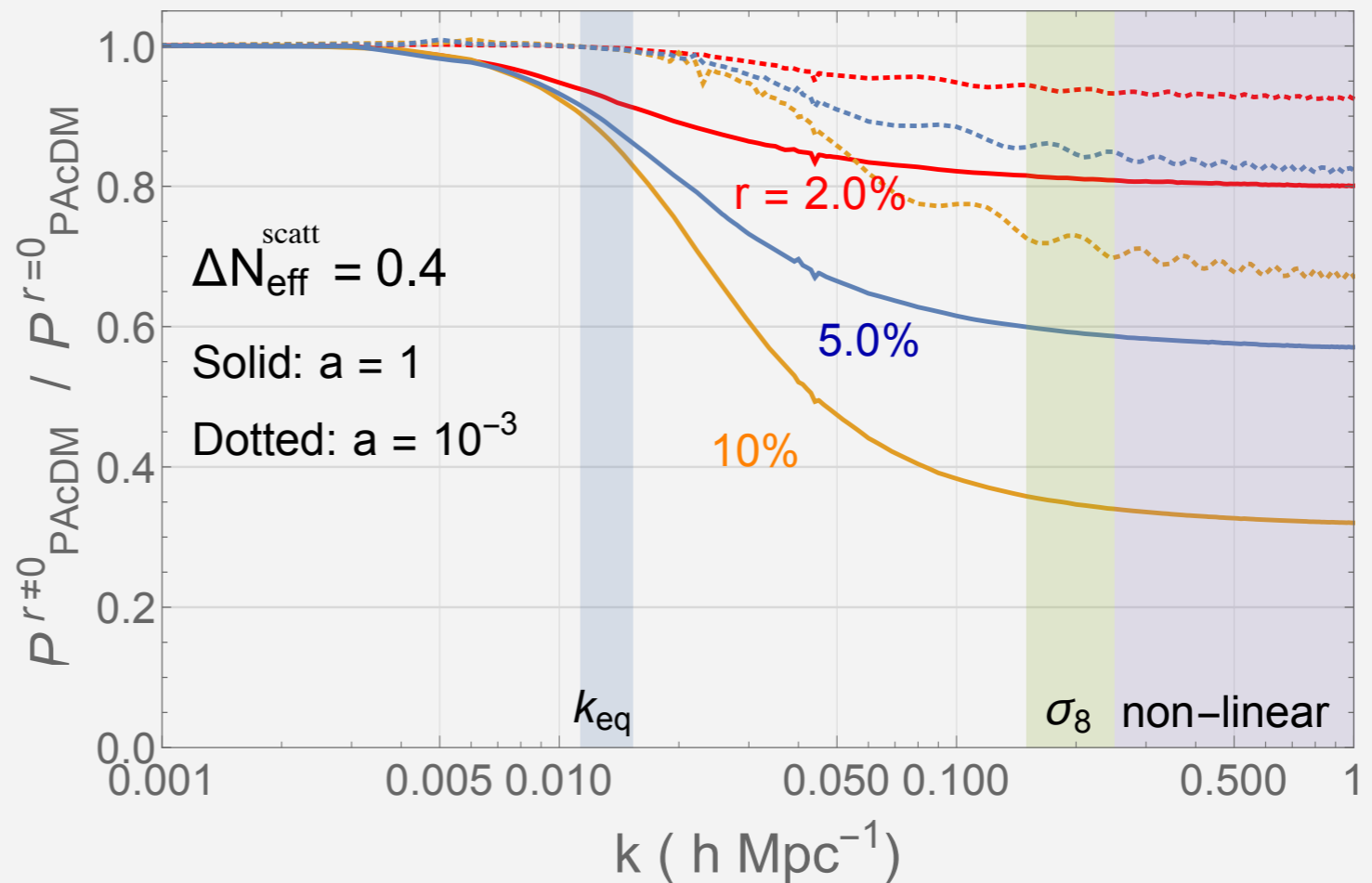
Allows an analytical study

$$\frac{P(r)}{P(0)} \simeq (1 - 2r) \left(\frac{a}{a_{\text{eq}}} \right)^{-1.2r}$$

$$r \equiv \Omega_2 / \Omega_{\text{DM}}$$

Solving Sigma8 problem with PAcDM

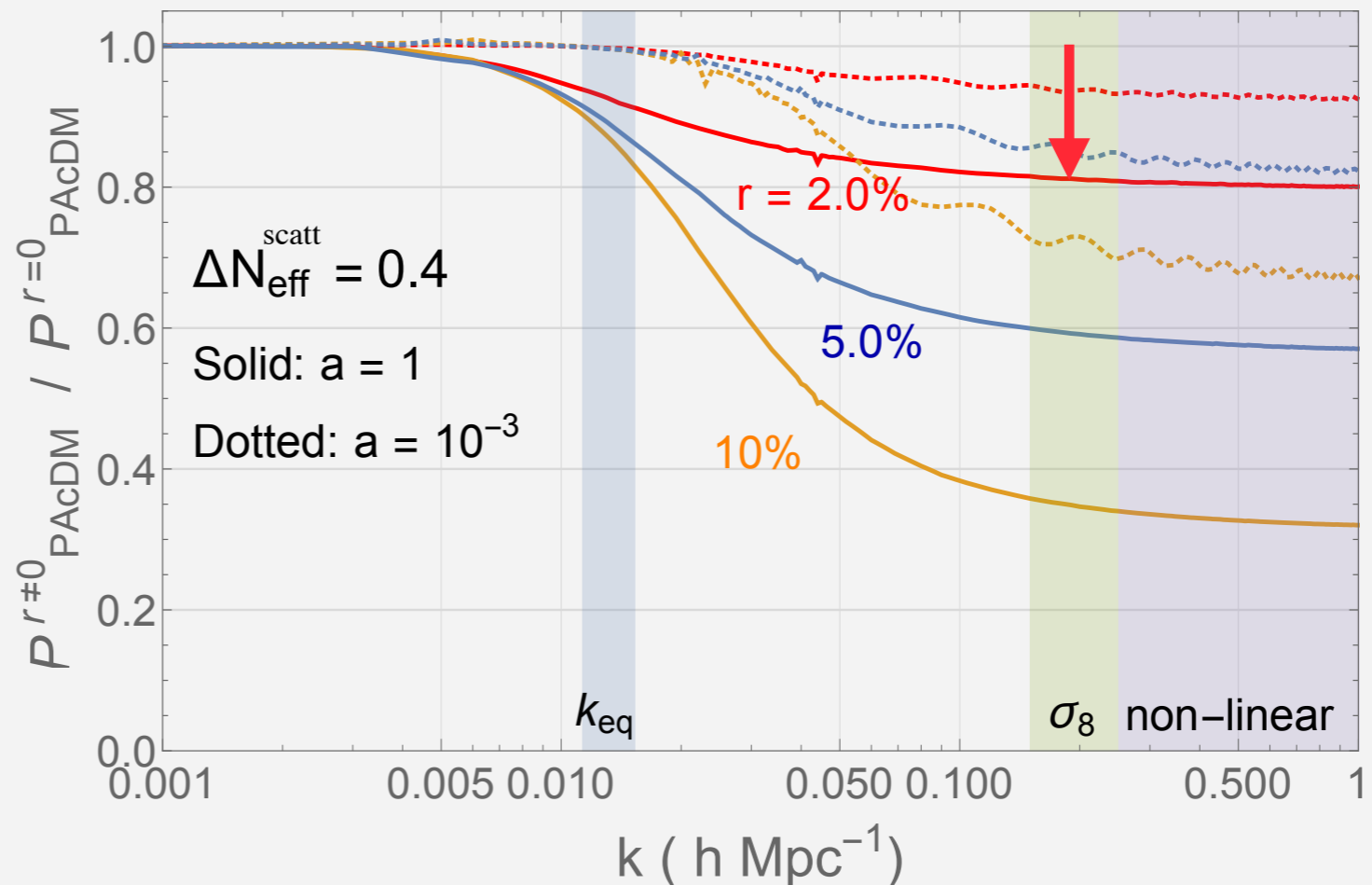
$$\left[\frac{\sigma_8(\Lambda\text{CDM})}{\sigma_8(\text{PAcDM})} \right]^2 \sim$$



$$r \equiv \Omega_2 / \Omega_{\text{DM}}$$

Solving Sigma8 problem with PAcDM

$$\left[\frac{\sigma_8(\Lambda\text{CDM})}{\sigma_8(\text{PAcDM})} \right]^2 \sim$$



Need $\sim 2\%$ acoustic DM to solve the σ_8 problem

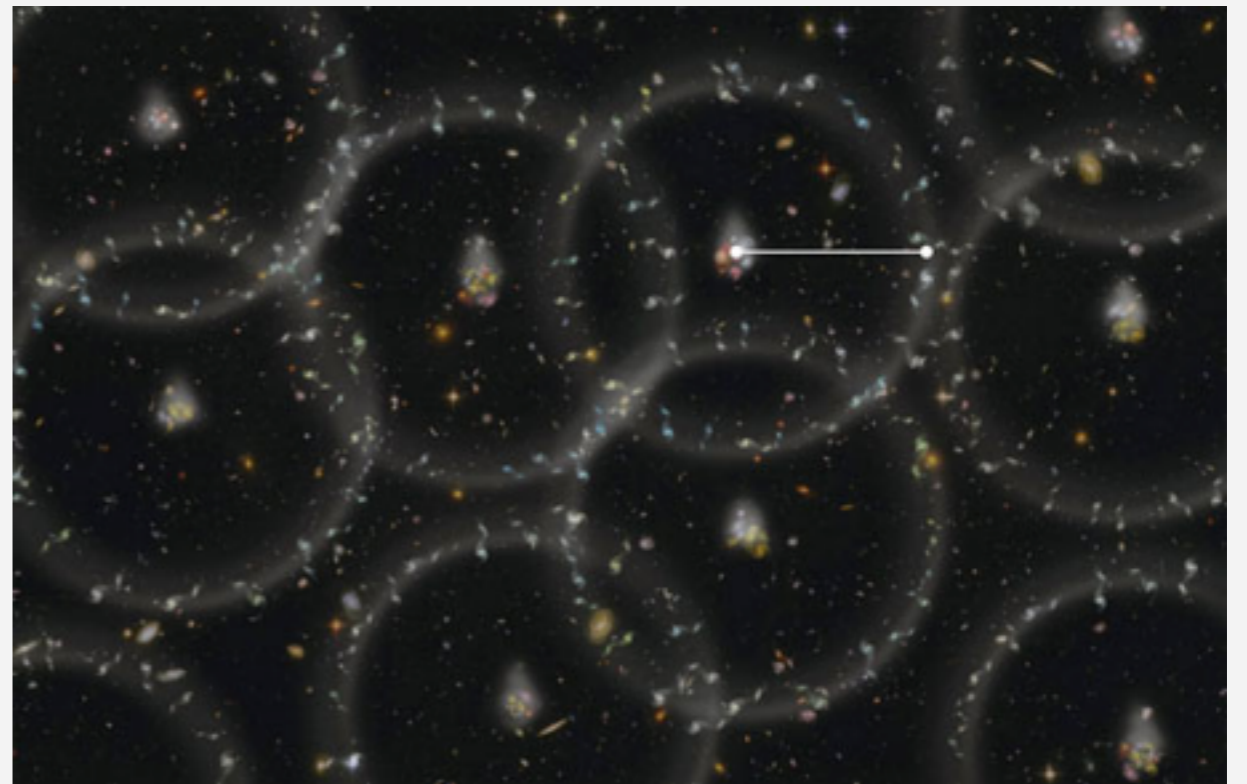
2% density is easy to obtained

When both DM particles are WIMP-like and having thermal freeze out through a heavy mediator

$$\frac{\Omega_2}{\Omega_1} \simeq \left(\frac{m_2}{m_1} \right)^2$$

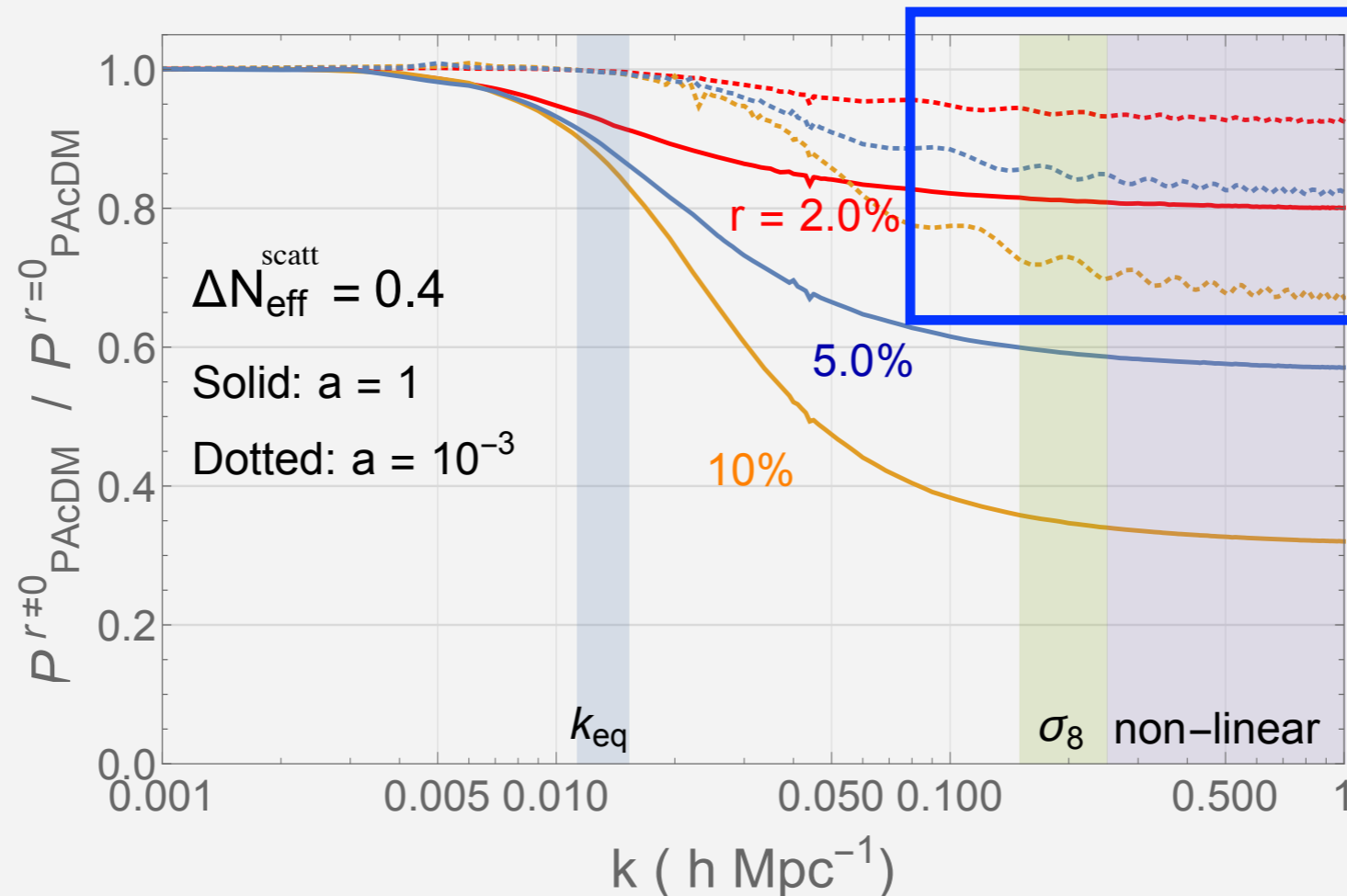
Only need $m_1 \simeq 7 m_2$ to obtain the 2% ratio
(assuming equal couplings)

Slowing Down the Structure Formation



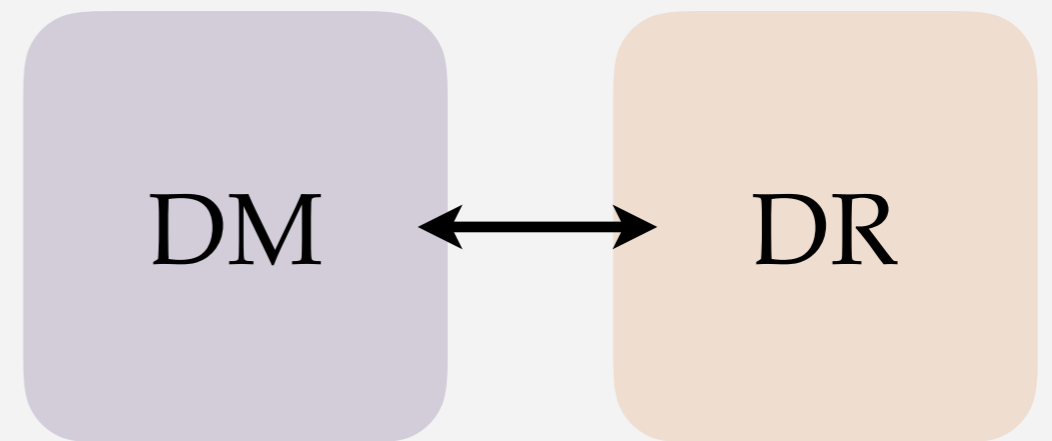
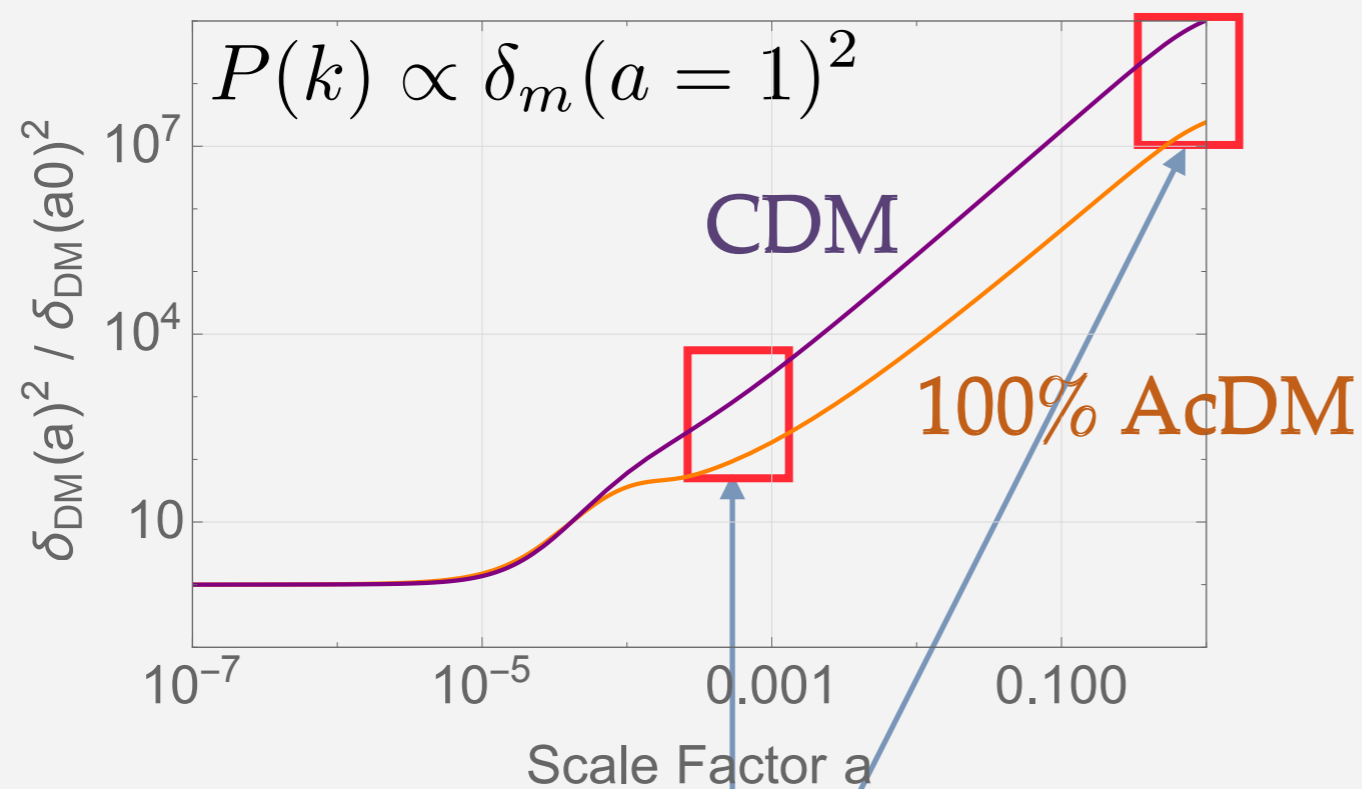
A cartoon produced by the BOSS project showing the spheres of baryons around the initial dark matter clumps

Smaller suppression at the CMB time



Correction to the power spectrum is smaller during the CMB time. Why?

In the **Quasi-Acoustic Oscillation** case

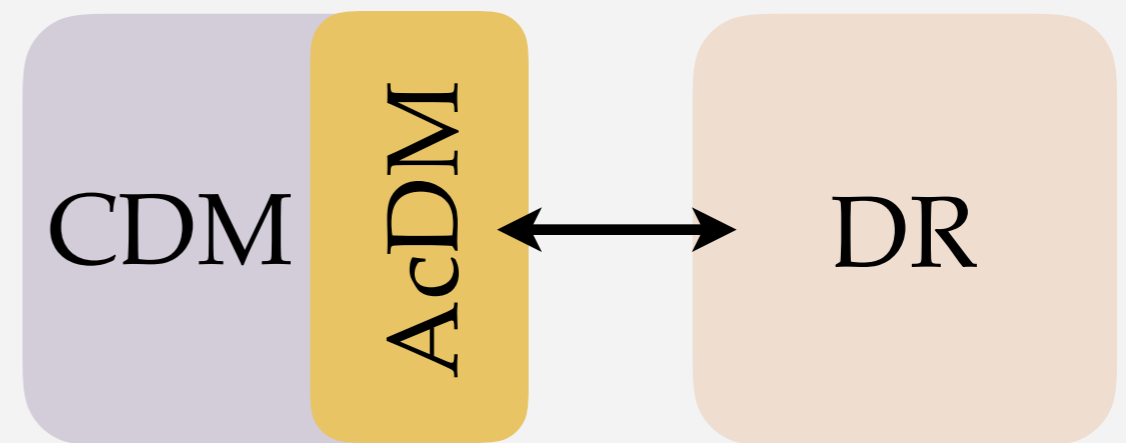
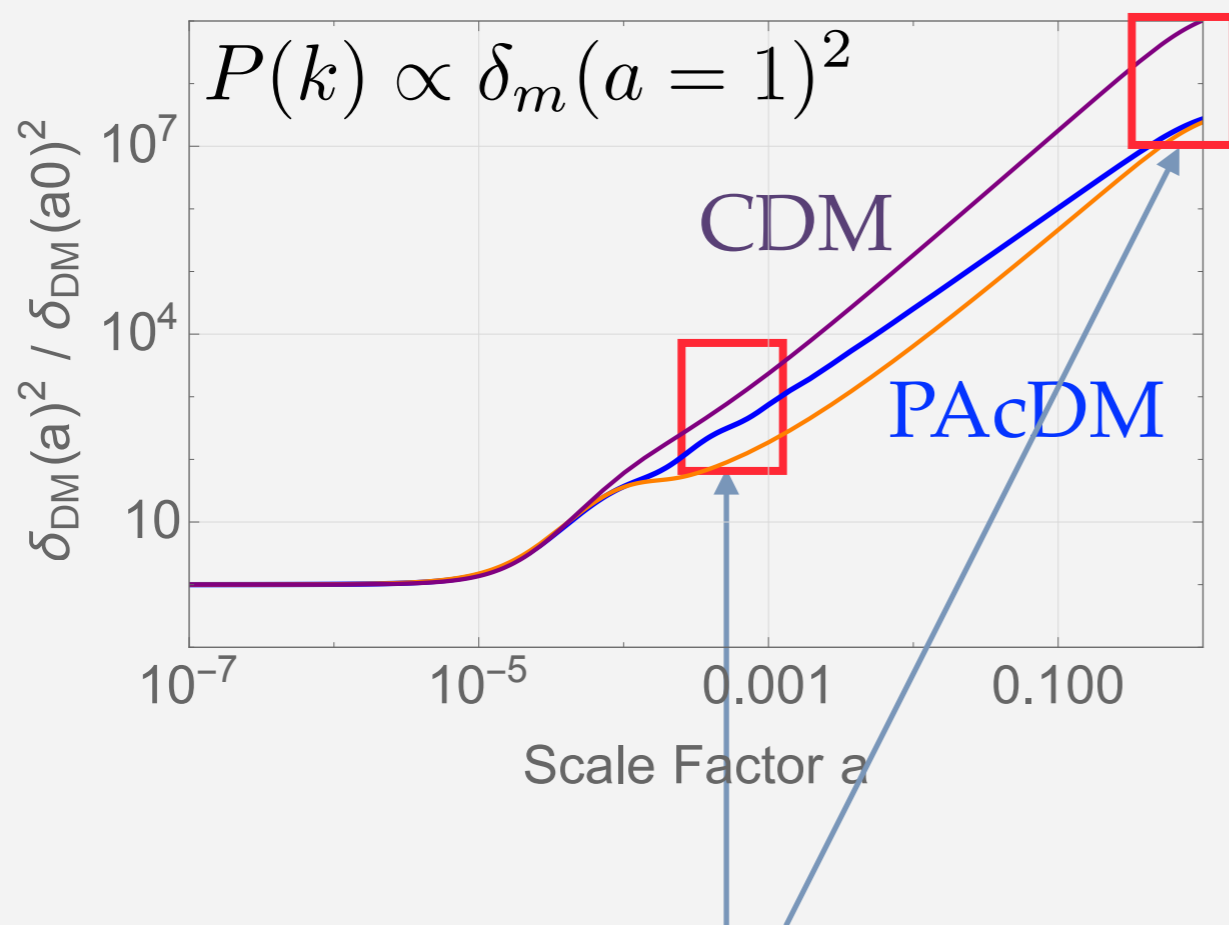


For TeV scale DM

$$\hat{\alpha} \sim 10^{-8.5}$$

Similar damping between today / CMB time

In the Partially-Acoustic Oscillation case

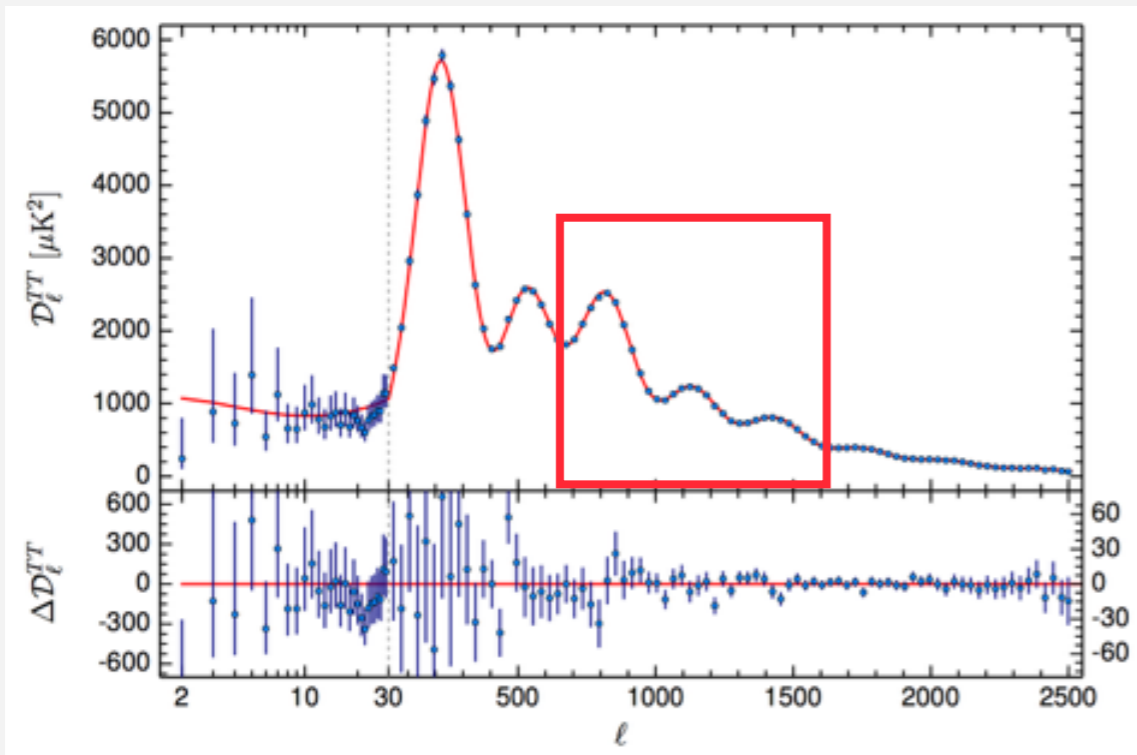


$$\frac{P(r)}{P(0)} \simeq (1 - 2r) \left(\frac{a}{a_{eq}} \right)^{-1.2r}$$

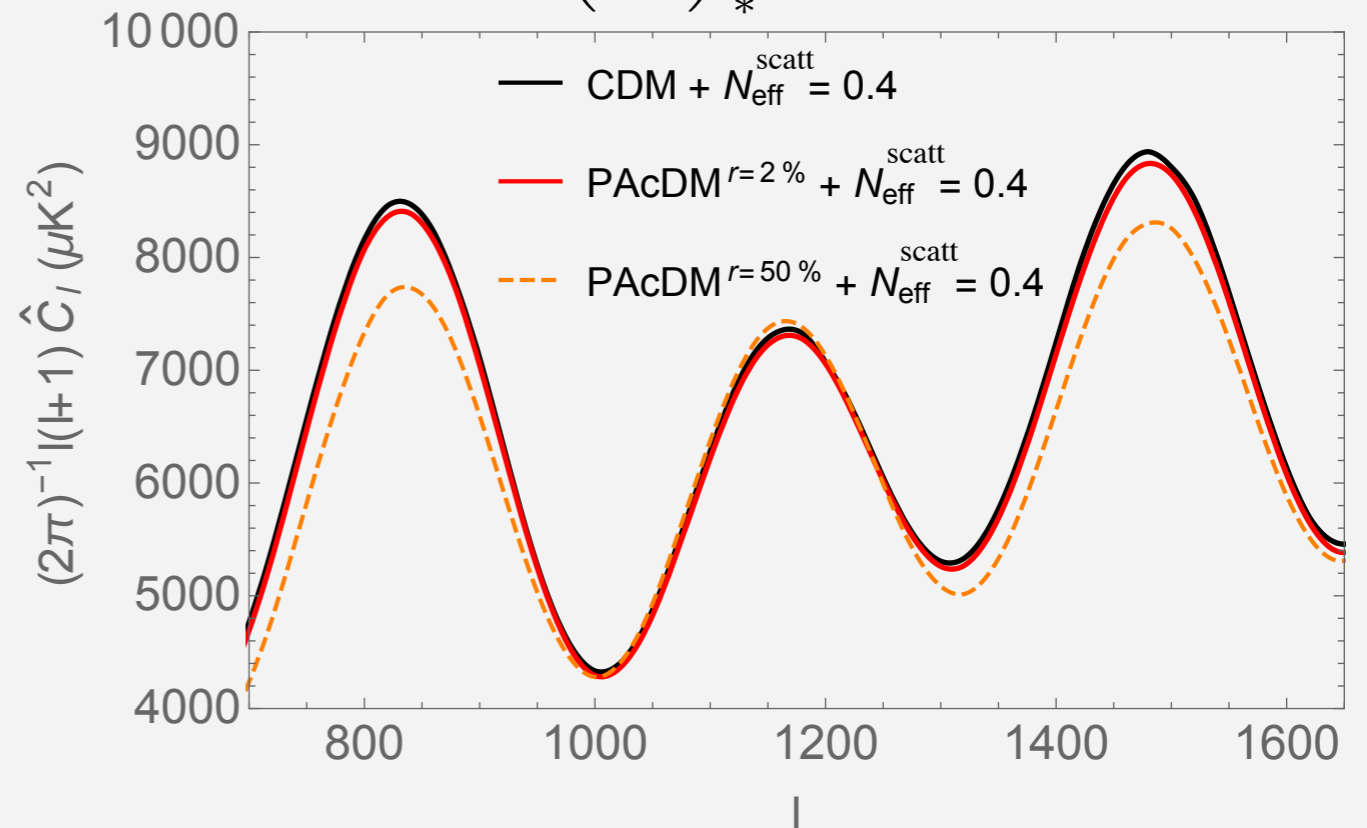
Structure grows **slower** comparing to CDM
 Smaller correction to the CMB spectrum

Correction to the CMB spectrum

Planck 1502.01589



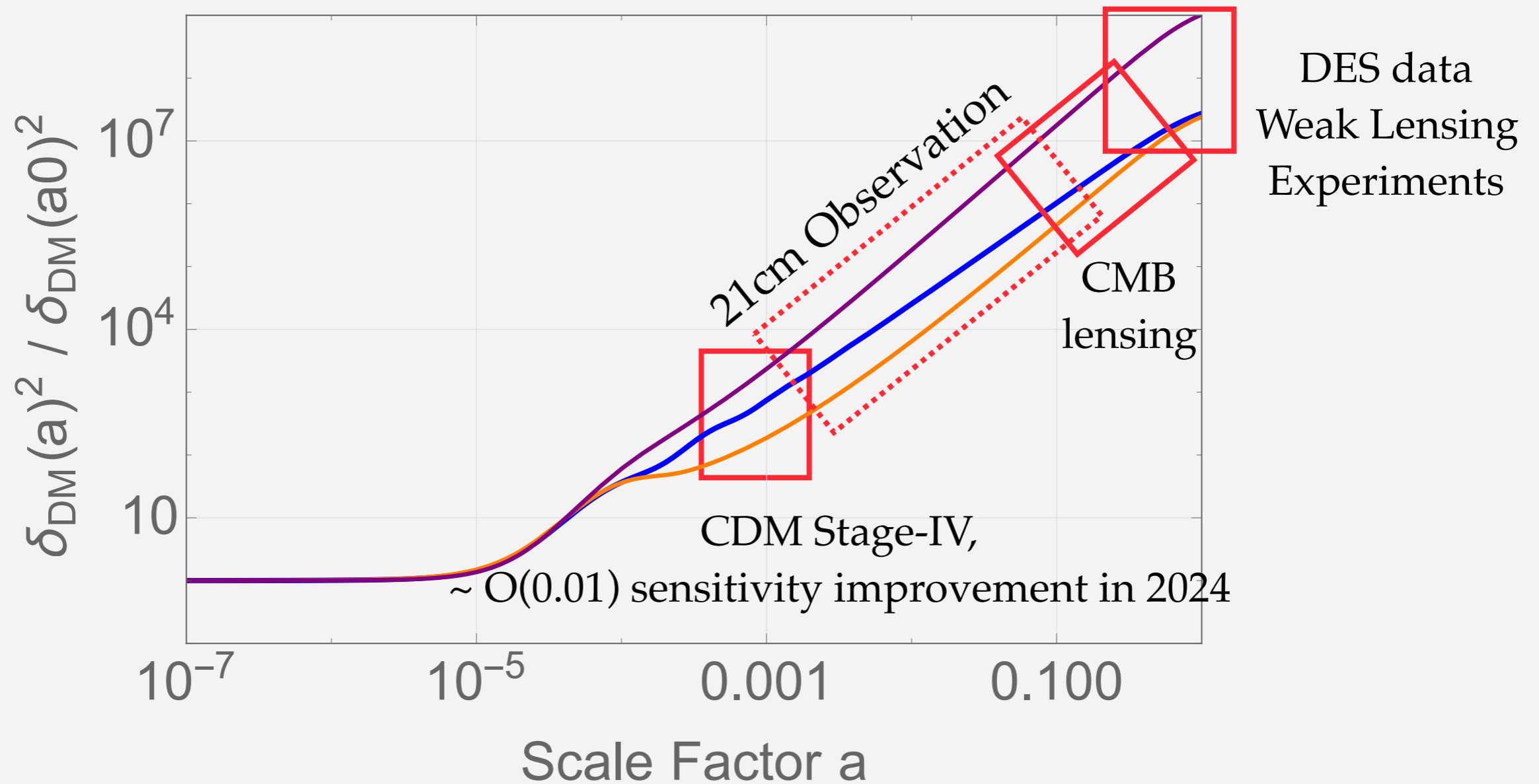
$$\left(\frac{\delta T}{T}\right)_* \equiv \frac{\delta\gamma}{4} + \psi$$



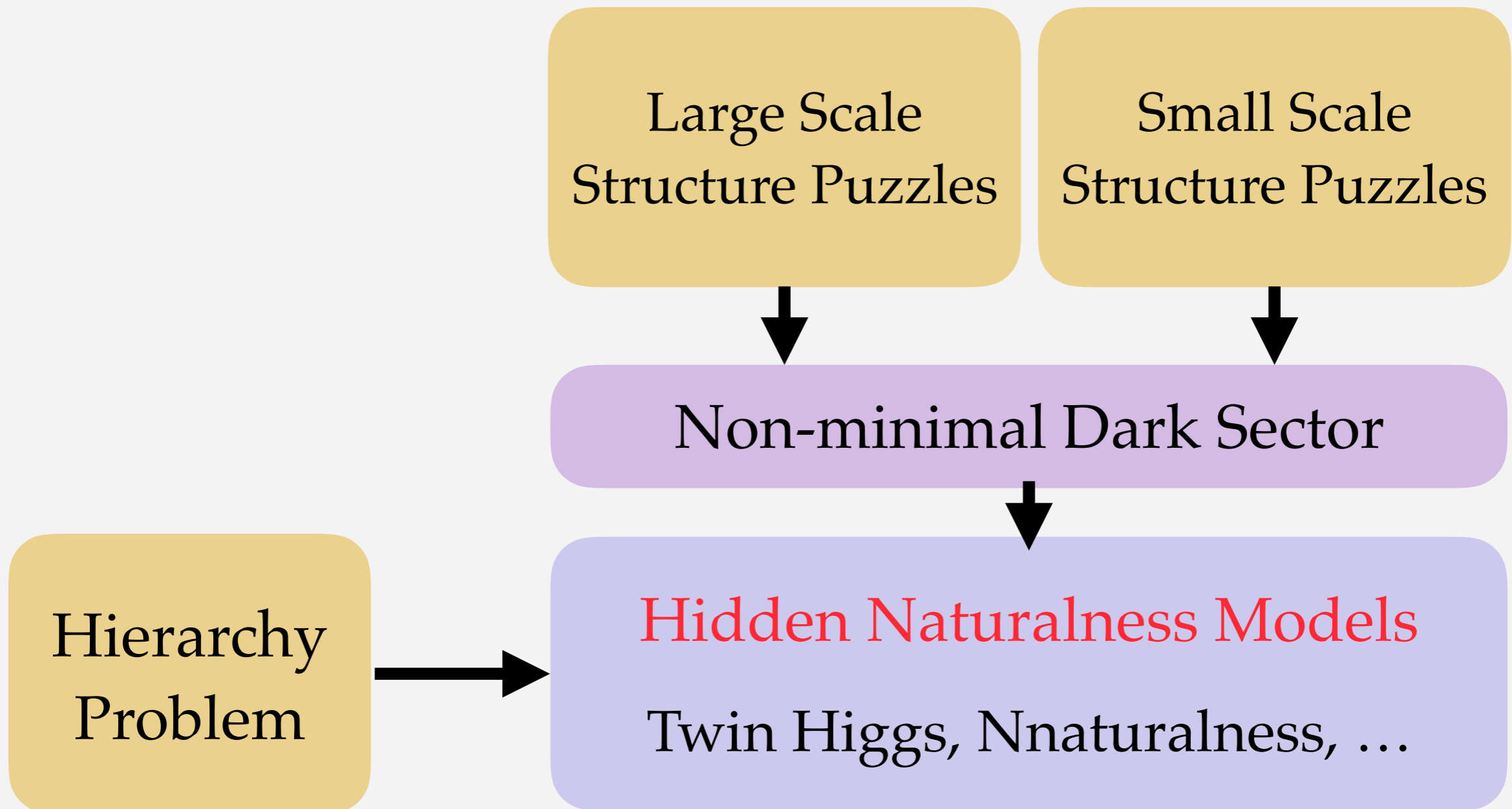
The pressure from dark fluid suppresses the compression peaks and enhances the expansion peaks

When $r = 2\%$, the correction to CMB is less than $\sim 2\%$, smaller than $> 5\%$ error bar in Planck result

Clear answer from future experiments



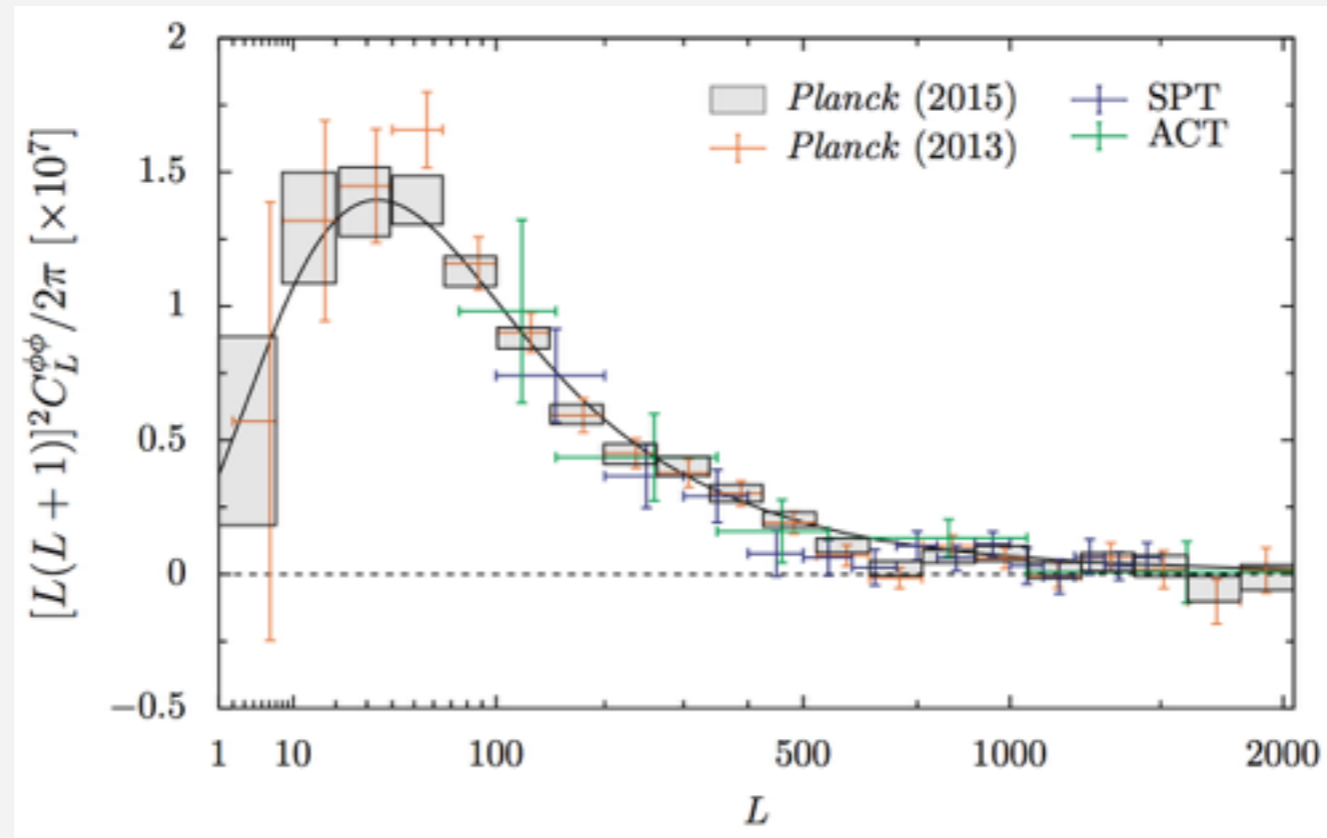
A clue to even deeper physics?



arXiv: 1611.05879, Valentina Prilepina and YT

CMB Lensing

Planck. 1502.01591



$$\ell^4 C_\ell^{\phi\phi} \propto \int_0^{\chi_*} d\chi (1 - \chi/\chi_*)^2 k\phi^2(k, a) g^2(\chi)$$

The smallest error bar (Planck) is 5% at $L \sim 150$
PAcDM gives a $\sim 2.5\%$ correction when $r = 2\%$

Conclusion

Large Scale Structure is sensitive to the dark sector dynamics

Acoustic Dark Oscillation

suppresses the matter power spectrum

A smaller ratio of Cold DM

change the power-law growth of matter density spectrum

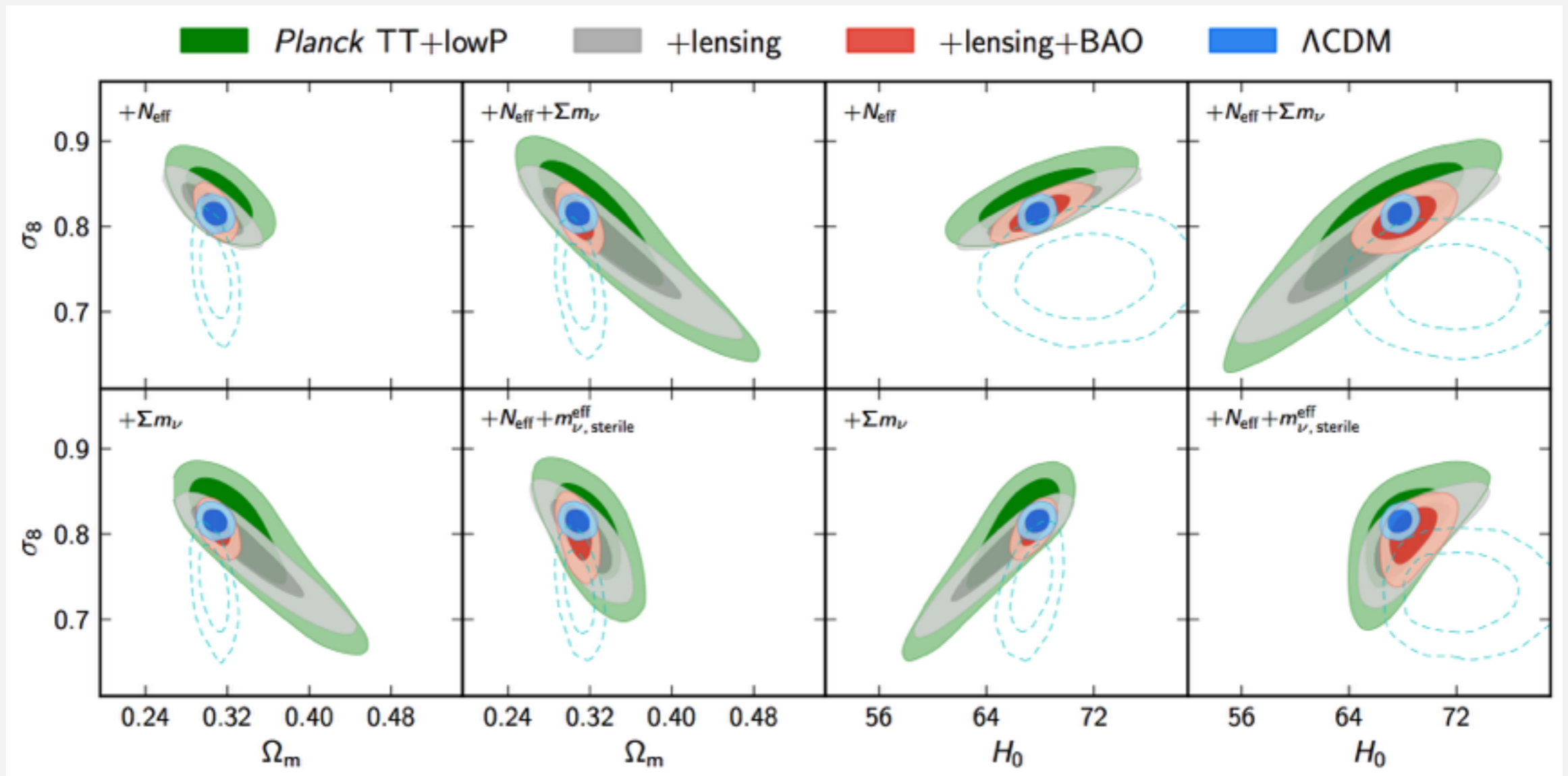
Having Dark Radiation

change the expansion, different effects on CMB between
free-streaming / self-scattering

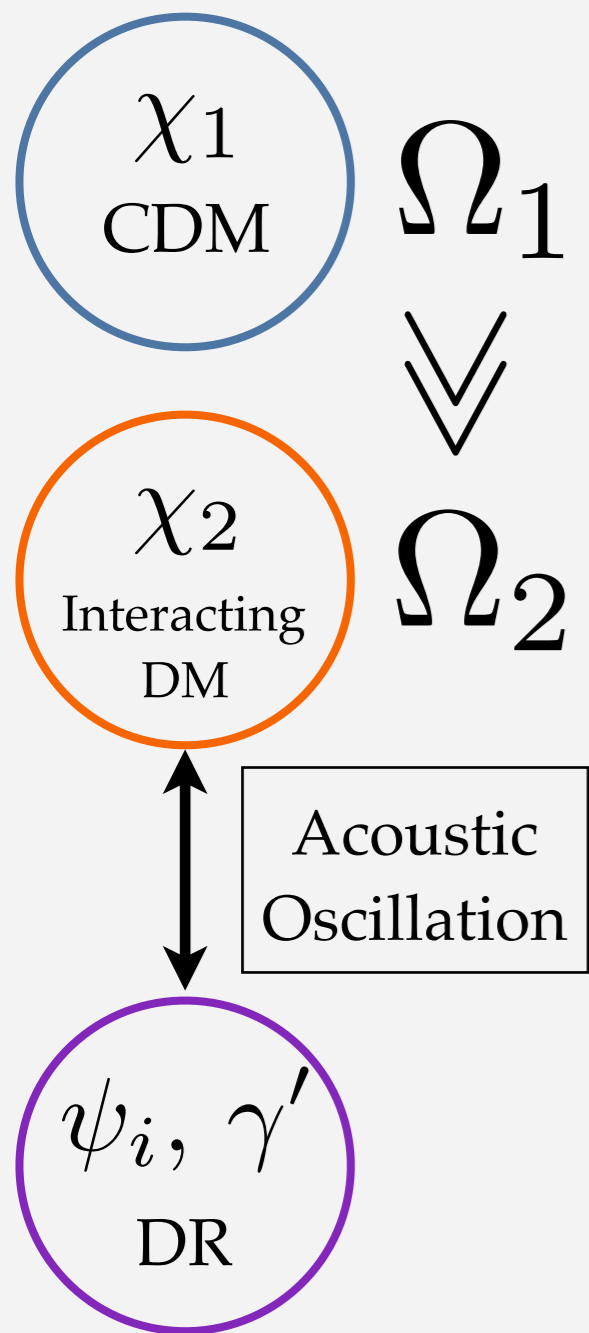
May also change the small scale structure

Working on it now, stay tuned!

Backup Slides



In the partially acoustic case



Acoustic
Oscillation

$$\Rightarrow \delta_1 \gg \delta_2$$

DM density contrast is determined by χ_1

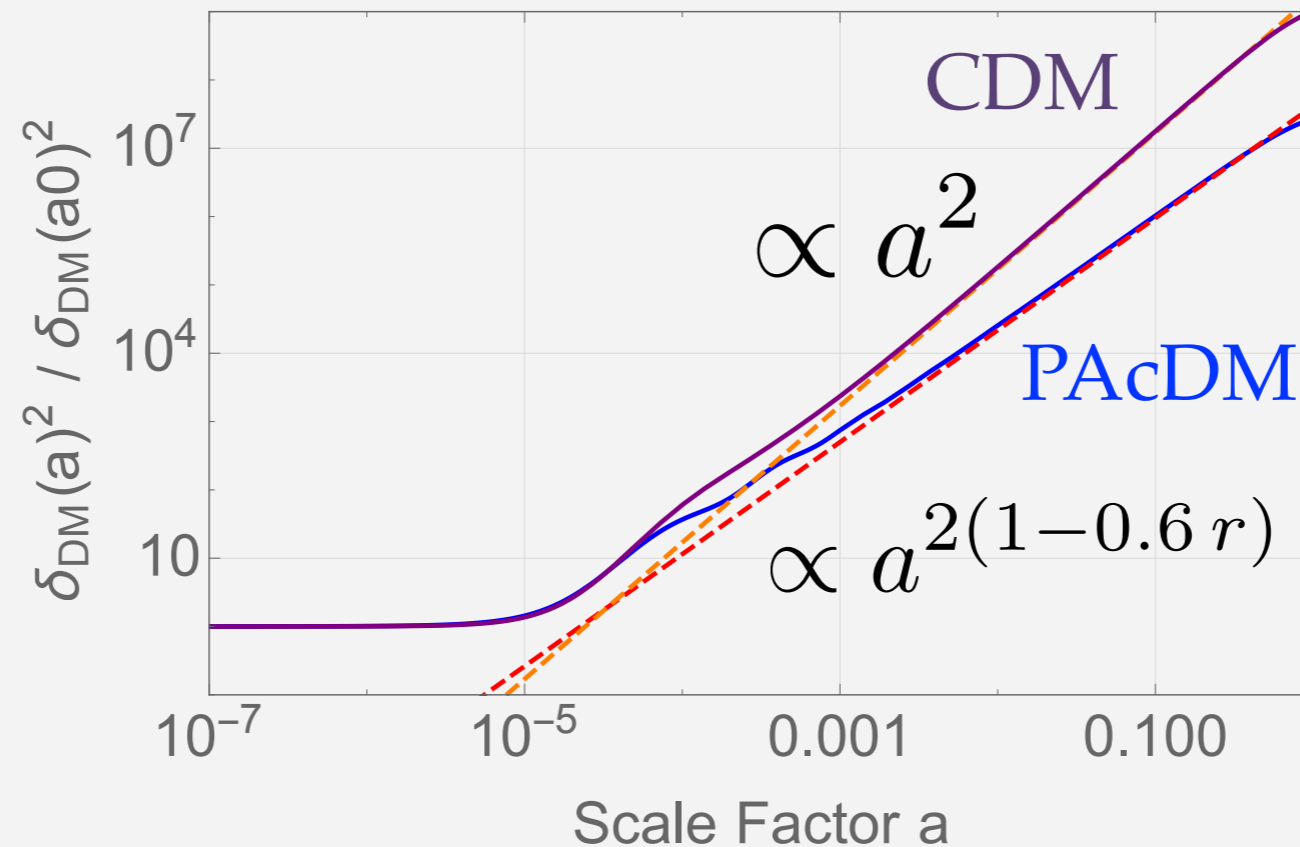
$$\ddot{\delta}_1 + \frac{2}{\tau} \dot{\delta}_1 = -k^2 \phi$$

$$k^2 \phi \simeq -4\pi G a^2 (\delta_1 \rho_1 + \delta_2 \rho_2)$$

$$= -\frac{6}{\tau} (1 - r) \delta_1 \quad r \equiv \frac{\rho_2}{\rho_{DM}}$$

The < 100% CDM case

$$\ddot{\delta}_1 + \frac{2}{\tau} \dot{\delta}_1 = \frac{6}{\tau} (1 - r) \delta_1 \quad \Rightarrow \quad \delta_1 \propto \left(\frac{a}{a_{eq}} \right)^{1 - 0.6 r + \mathcal{O}(r^2)}$$



Boltzmann Equation in Conformal Newtonian Gauge

$$\dot{\delta}_D = -\theta_D + 3\dot{\psi} \quad \left(\frac{d\rho}{dt} = -\rho \nabla \cdot \vec{v} \right)$$

$$\dot{\theta}_D = -\frac{\dot{a}}{a}\theta_D + k^2\psi + \underline{a\Gamma(\theta_R - \theta_D)} \quad \Gamma \equiv \frac{1}{\langle p_D^2 \rangle} \frac{d\langle \delta p_D^2 \rangle}{dt}$$

$$\dot{\delta}_R = -\frac{4}{3}\theta_R + 4\dot{\psi}$$

$$\dot{\theta}_R = \frac{k^2}{4}\delta_R + k^2\psi + \underline{Ra\Gamma(\theta_D - \theta_R)}$$

$$\theta_s \equiv \partial_i v_s^i$$

Velocity Divergent

Metric Perturbation

$$ds^2 = a^2(\tau) [-(1 + 2\psi)d\tau^2 + (1 - 2\phi)\delta_{ij}dx^i dx^j]$$

No free-streaming particle => $\phi = \psi$

Boltzmann Equation in Conformal Newtonian Gauge

$$\dot{\delta}_D = -\theta_D + 3\dot{\psi} \quad \left(\frac{d\rho}{dt} = -\rho \nabla \cdot \vec{v} \right)$$

$$\dot{\theta}_D = -\frac{\dot{a}}{a}\theta_D + k^2\psi + \underline{a\Gamma(\theta_R - \theta_D)} \quad \Gamma \equiv \frac{1}{\langle p_D^2 \rangle} \frac{d\langle \delta p_D^2 \rangle}{dt}$$

$$\dot{\delta}_R = -\frac{4}{3}\theta_R + 4\dot{\psi}$$

$$\dot{\theta}_R = \frac{k^2}{4}\delta_R + k^2\psi + \underline{Ra\Gamma(\theta_D - \theta_R)}$$

Tightly coupled DM-DR (similar to the baryon-photon system):

$$\Gamma \gg H \quad \Rightarrow \quad a\Gamma \gg \tau^{-1}$$